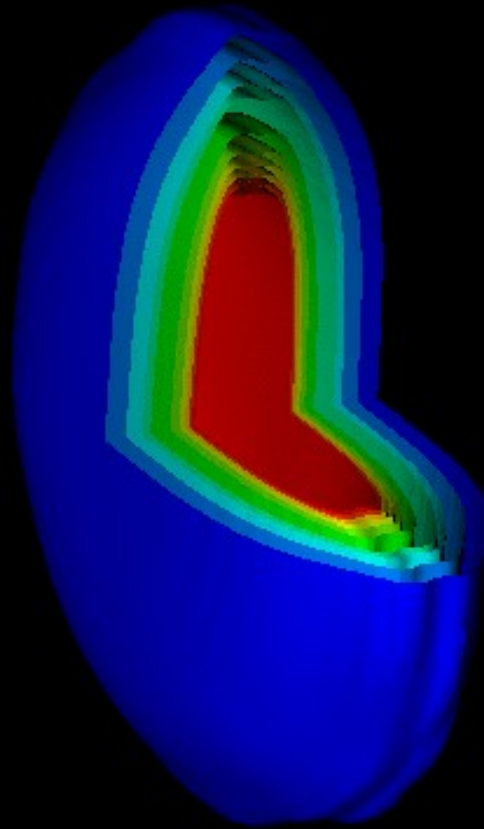
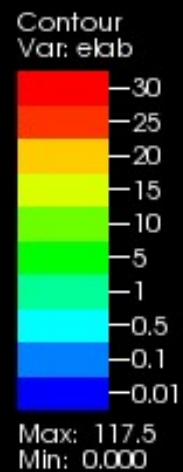
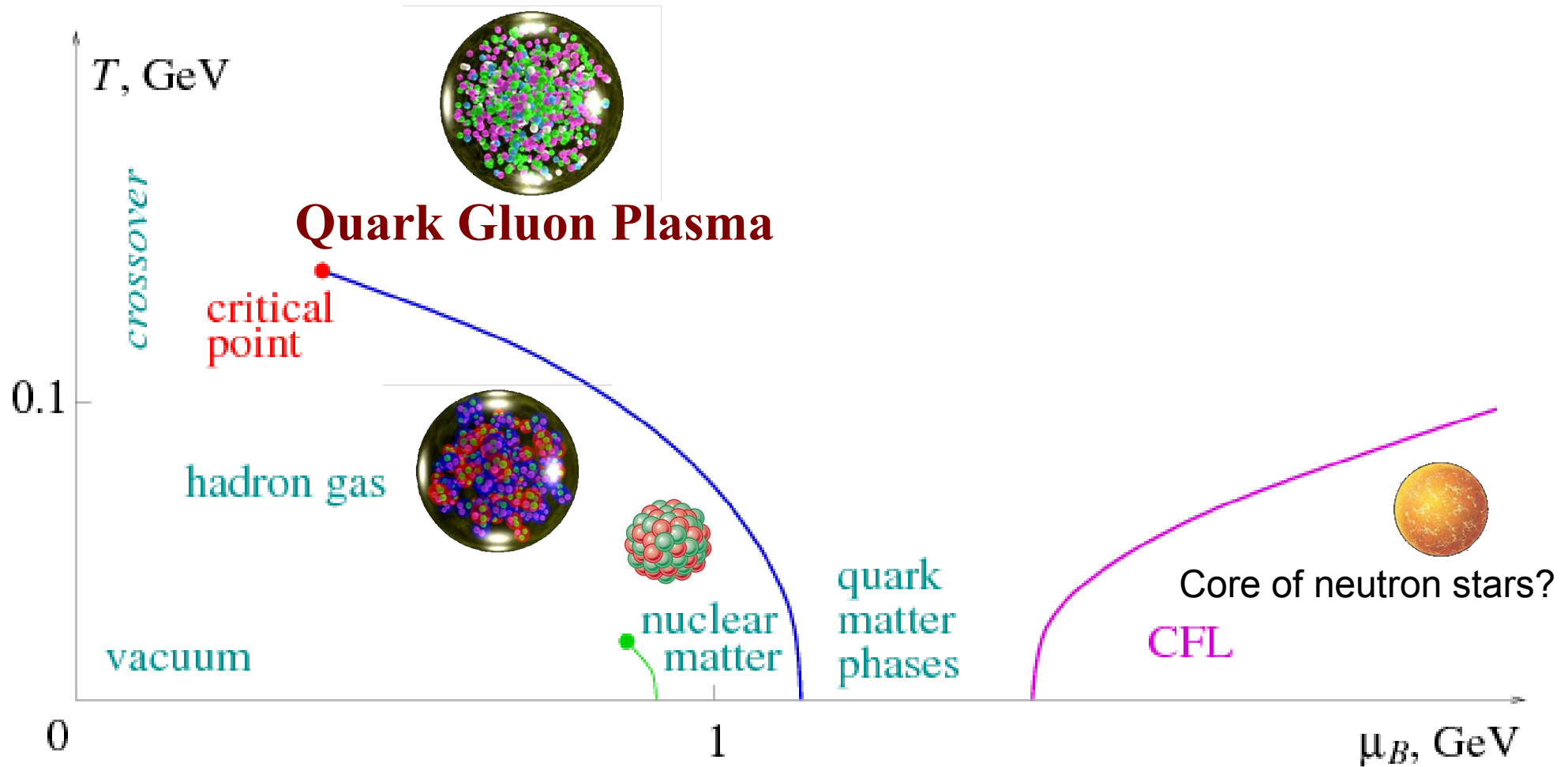


Transverse energy: measuring the energy density of the QGP



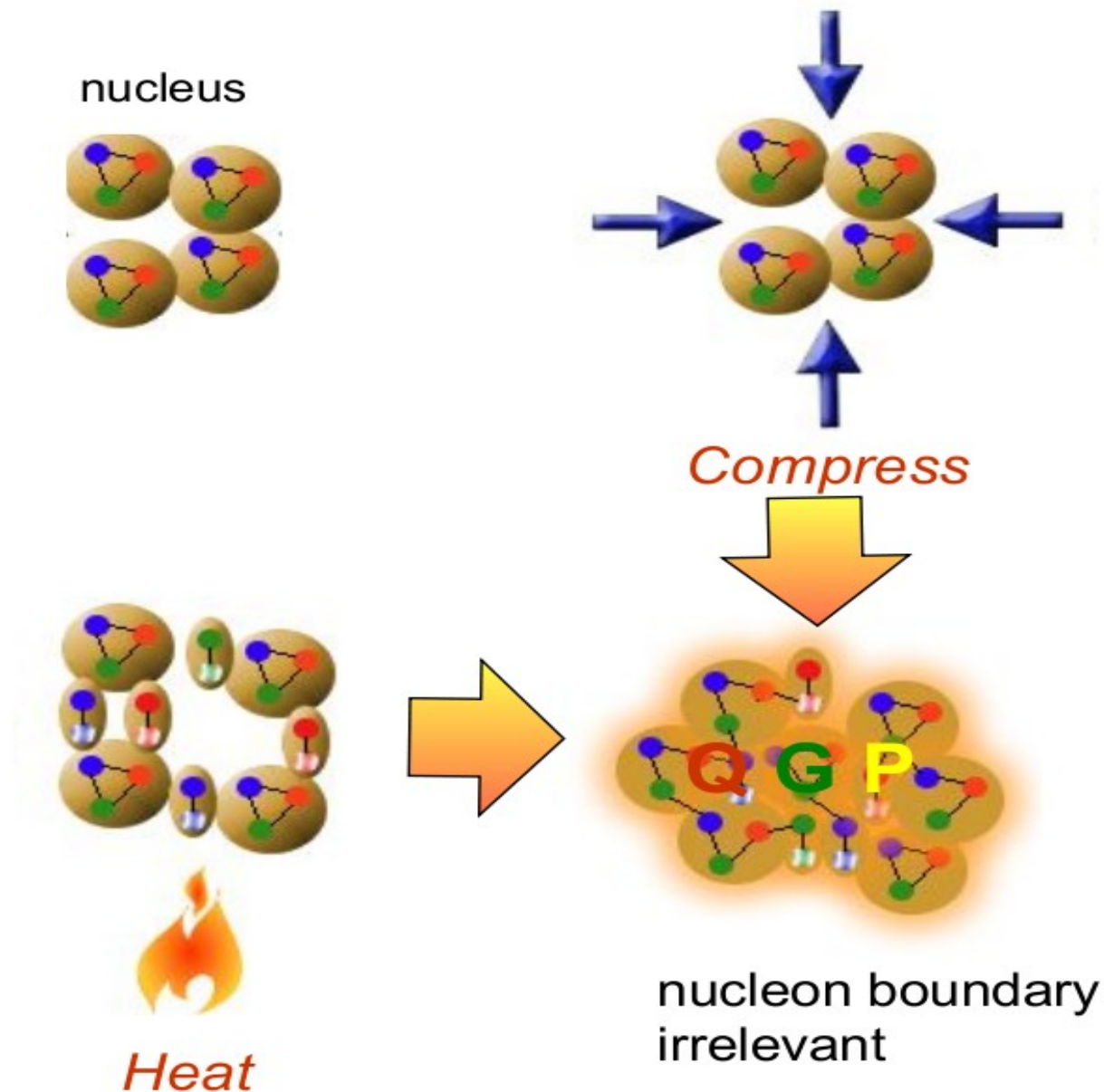
Christine Nattrass
University of Tennessee at Knoxville

Phase diagram of nuclear matter

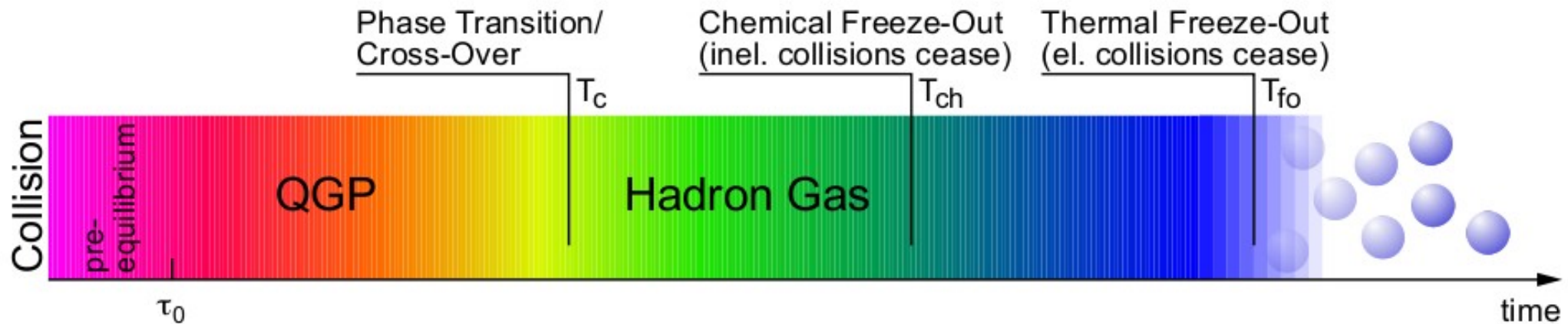
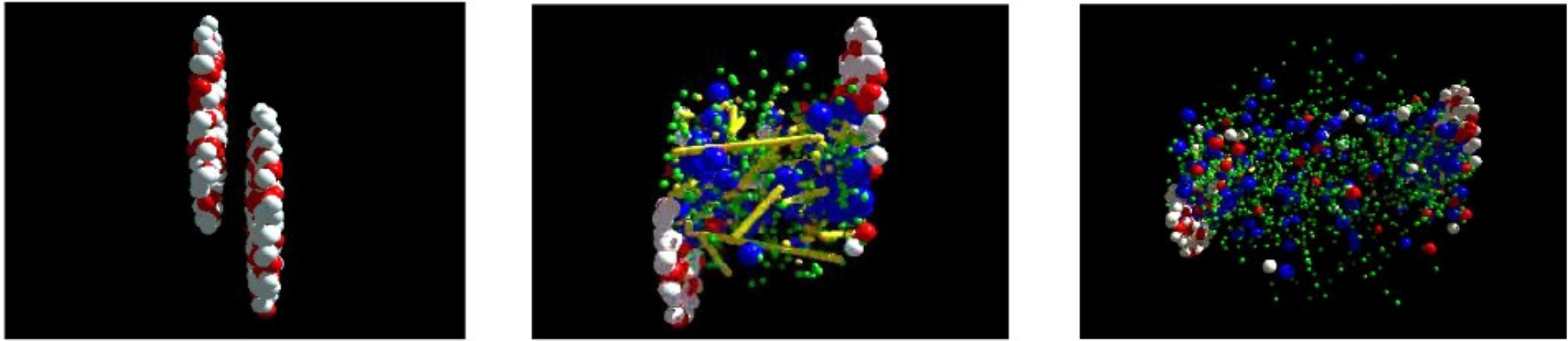


Quark Gluon Plasma – a *liquid* of quarks and gluons created at temperatures above ~ 170 MeV ($2 \cdot 10^{12}$ K) – over a million times hotter than the core of the sun

How to make a Quark Gluon Plasma



The phase transition in the laboratory



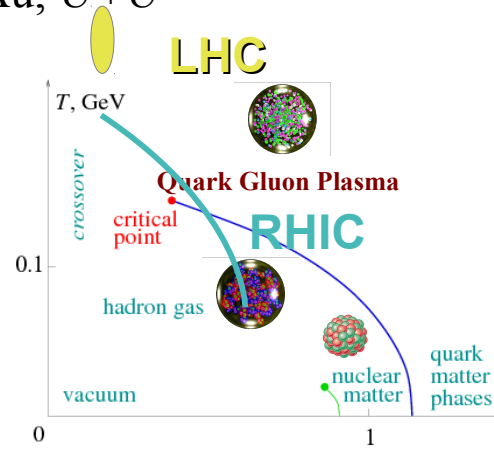
Relativistic Heavy Ion Collider



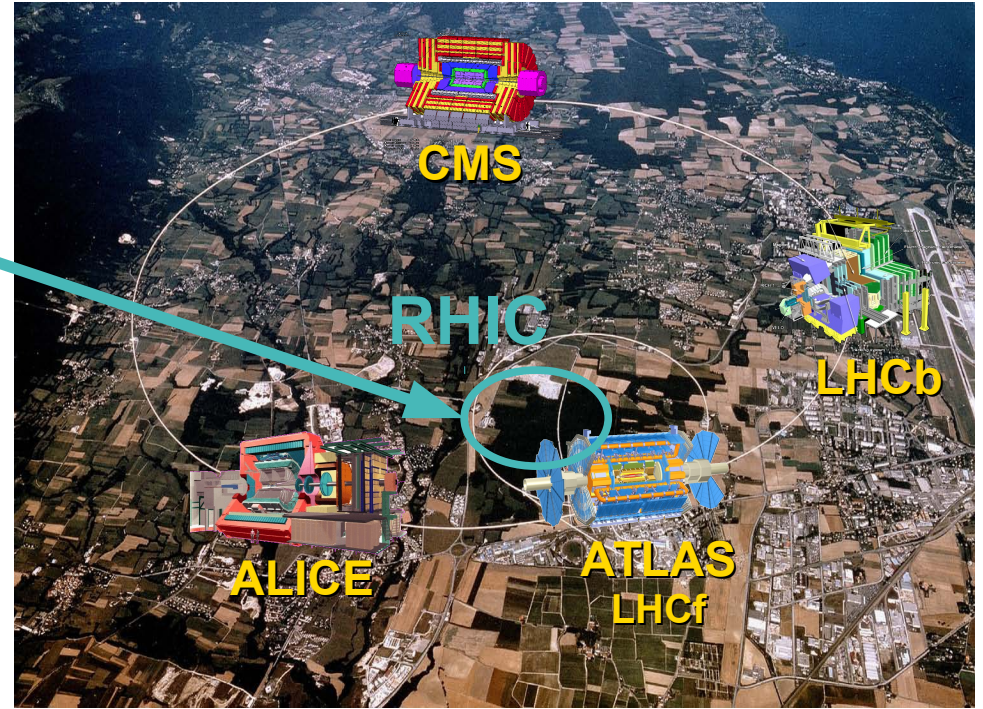
Upton, NY

1.2km diameter

p+p, d+Au, Cu+Cu, Au+Au, U+U

$$\sqrt{s}_{\text{NN}} = 9 - 200 \text{ GeV}$$


Large Hadron Collider

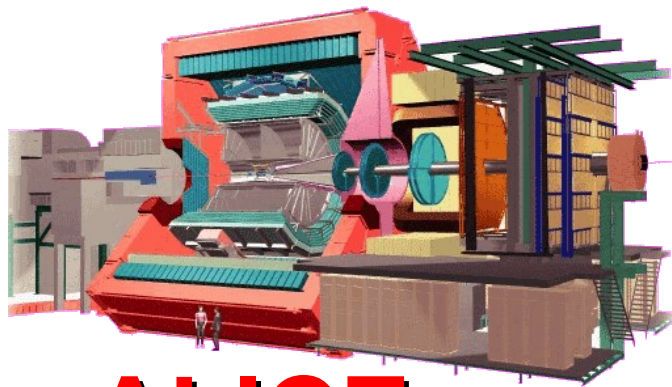


Geneva, Switzerland

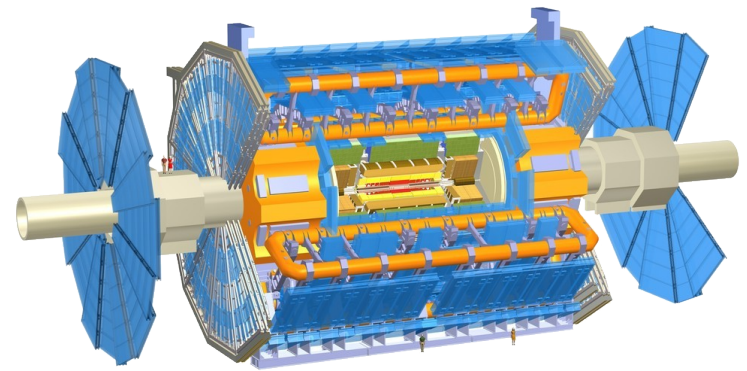
8.6km diameter

p+p, $p+Pb$, Pb+Pb

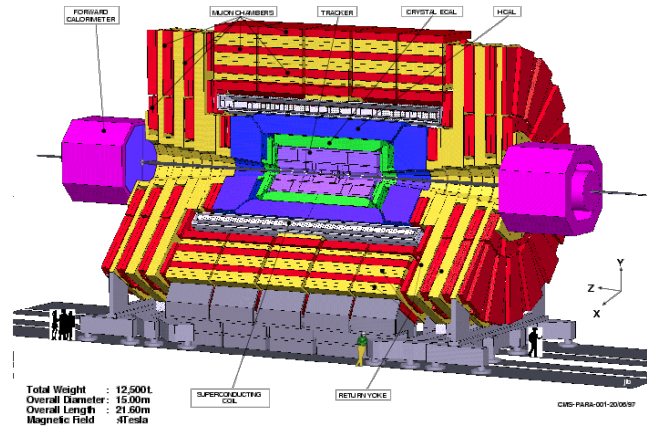
$$\sqrt{s}_{\text{NN}} = 2.76 \text{ GeV}, 5.5 \text{ TeV}$$



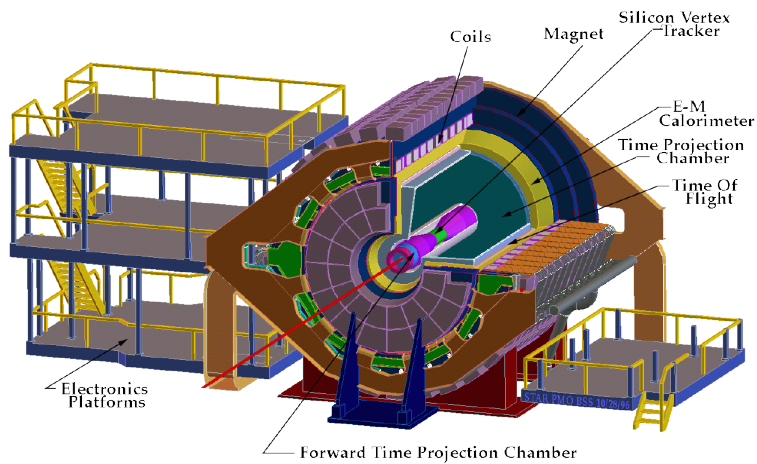
ALICE



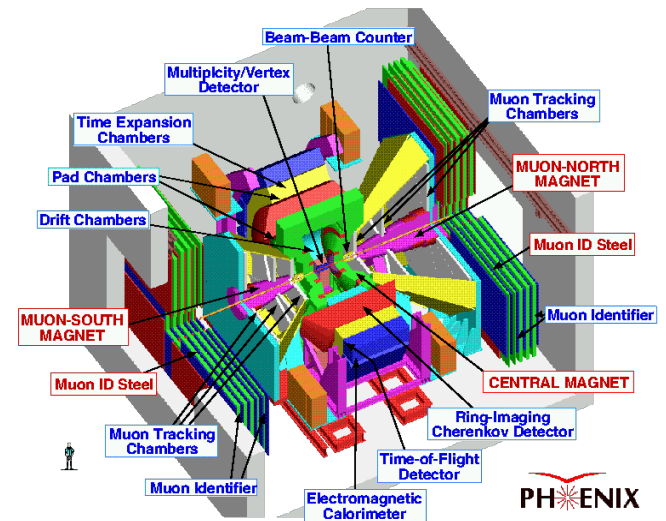
ATLAS



CMS

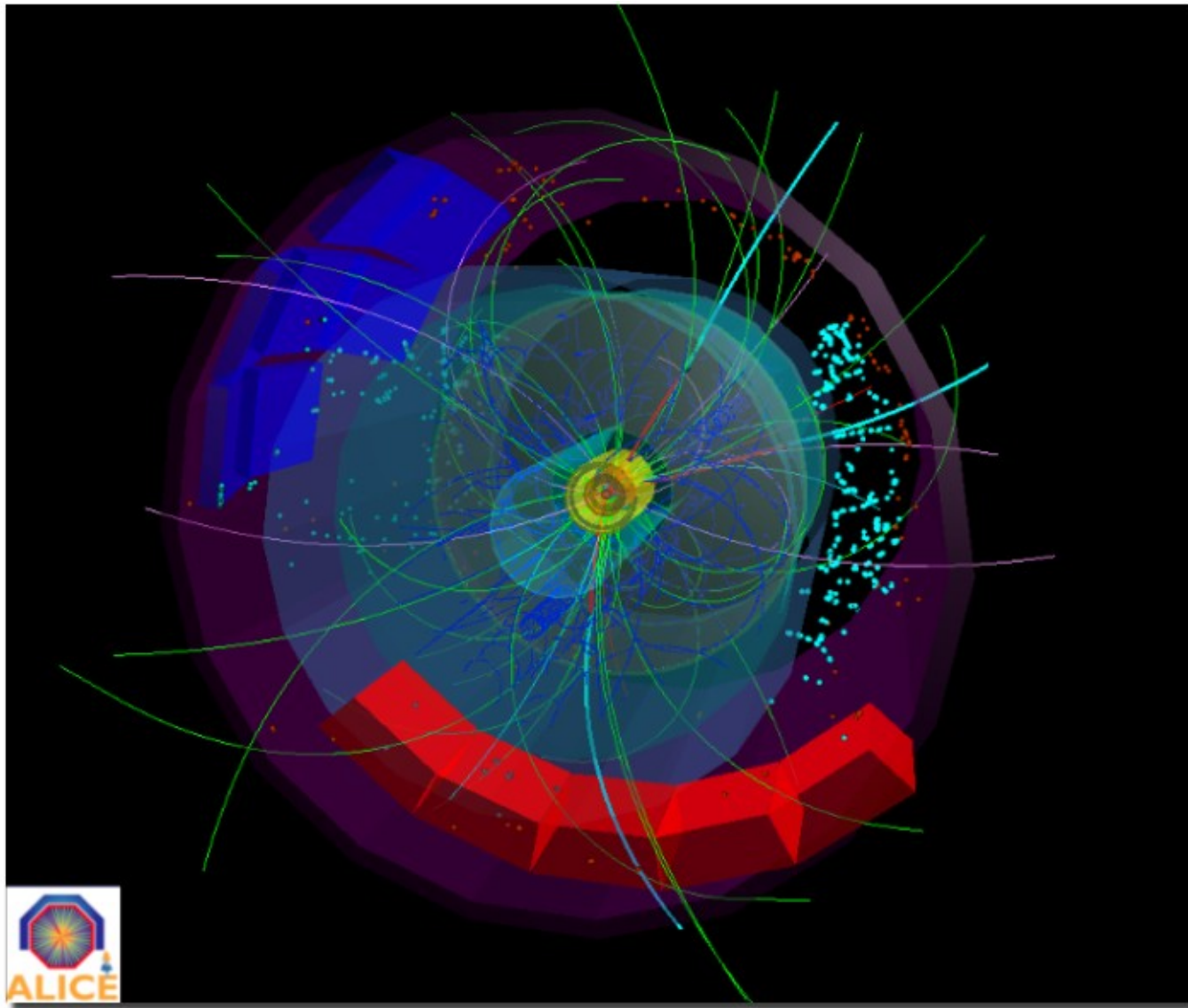


STAR



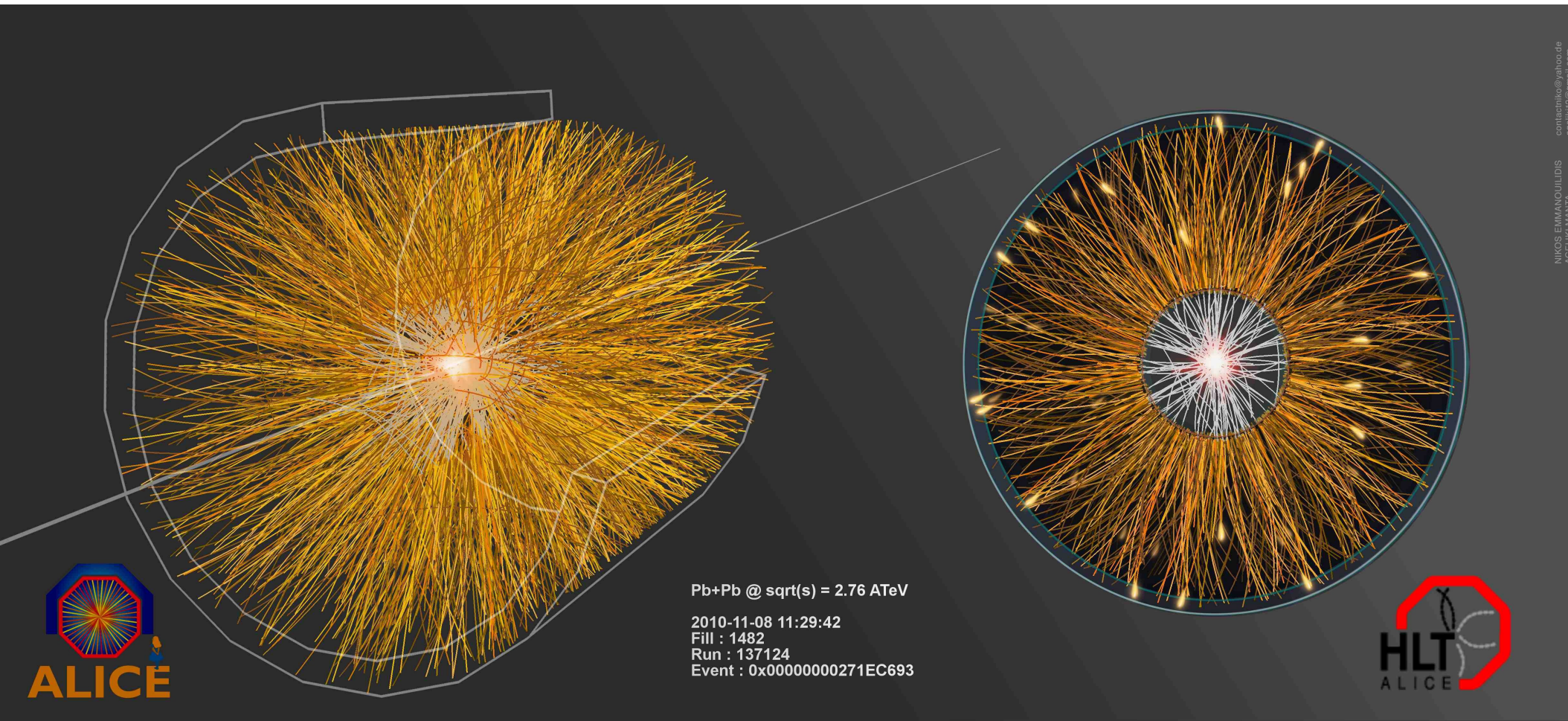
PHENIX

p+p collisions



3D image of each collision

Pb+Pb collisions



Comparison of colliders

	RHIC	LHC	
$\sqrt{s_{\text{NN}}} \text{ (GeV)}$	9-200	2760, 5500	<i>center of mass energy</i>
$dN_{\text{ch}}/d\eta$	~ 1200	~ 1600	<i>number of particles</i>
T/T_c	1.9	3.0-4.2	<i>temperature</i>
$\varepsilon \text{ (GeV/fm}^3\text{)}$	5	~ 15	<i>energy density</i>
$\tau_{\text{QGP}} \text{ (fm/c)}$	2-4	>10	<i>lifetime of QGP</i>

RHIC and LHC:

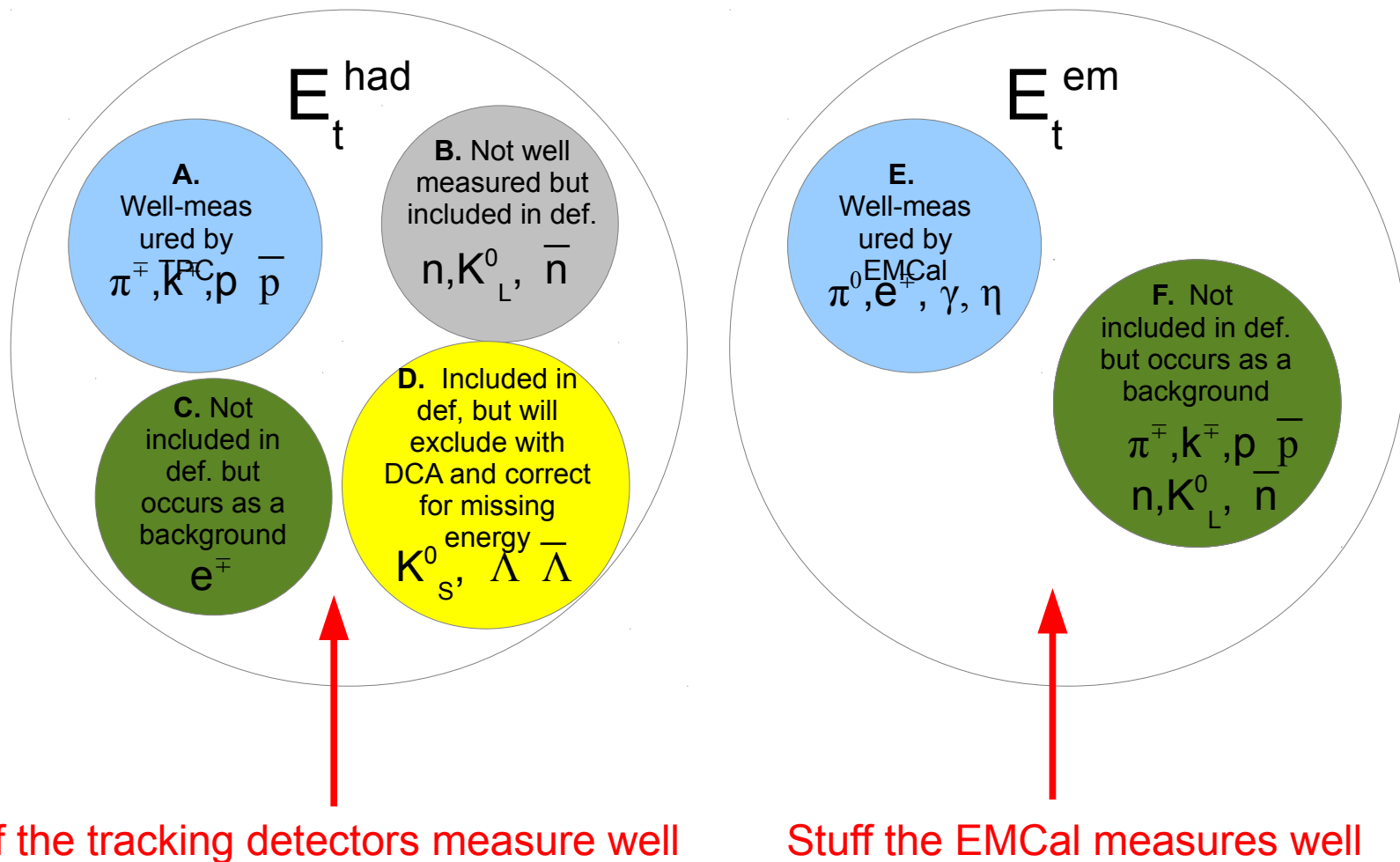
Cover 2 –3 decades of energy ($\sqrt{s_{\text{NN}}} = 9 \text{ GeV} - 5.5 \text{ TeV}$)

To discover the properties of hot nuclear matter at $T \sim 150 - 600 \text{ MeV}$

Measurements of transverse energy

- Energy directly transverse to the beam
 - $E_T = E \sin(\theta)$
- Methods
 - Calorimeter only
 - (Electromagnetic) calorimeter + tracking detectors
 - Tracking detectors only
 - Calculations from identified particle spectra

Hybrid method



$$E_T^{em}$$

Sum over
clusters

$$E_T^{em} = \frac{1}{f_{acc}} \frac{1}{f_{E_T min}} \left(\sum_i \delta_{matched} \frac{1}{\epsilon_\gamma f_{nonlin}} E_i \sin(\theta_i) - E_T^{kaons} - E_T^{ch.} - E_T^{(anti)neutrons} - E_T^{secondary} \right)$$

**Contributions
to final E_T^{em}
systematic
error**

Data driven

Input from data

$$\frac{1}{f_{acc}}$$

Geometric acceptance, not including dead channels

$$\frac{1}{f_{E_T min}}$$

Correction for minimum energy threshold **~6%**

$$\frac{1}{f_{nonlin}}$$

Correction for nonlinearity of detector response **~0.5%**

$$E_T^{kaons}$$

All energy deposited by K_S^0 , K_L^0 , K^\pm , including decays like $K_S^0 \rightarrow \pi^0 \pi^0 \rightarrow \gamma \gamma \gamma \gamma$ **<3%**

$$E_T^{ch}$$

Correction for other charged hadron deposits in calorimeter **~10-20%**

$$E_T^{(anti)neutrons}$$

Correction for (anti)neutron deposits in calorimeter **~1.5-5%**

$$E_T^{secondary}$$

Correction for deposits by particles from secondary interactions **<4 <5%**

$$\epsilon_\gamma$$

efficiency x acceptance within geometric acceptance of detector **~1%**

$$E_T^{had}$$

$$E_T^{had} = \frac{1}{f_{acc}} \frac{1}{f_{p_T cut}} \frac{1}{f_{neutral}} \sum_{i=0}^n f_{bg}^i(p_T) \frac{1}{f_{notID}} \frac{1}{eff(p_T^i)} E_i^{had} \sin(\theta^i)$$

$$\frac{1}{f_{acc}} \quad \text{Correction for the geometric acceptance – 1, with acceptance due to sector boundaries, etc. rolled into the track efficiency}$$

$$\frac{1}{f_{p_T cut}} \quad \text{Correction for the low } p_T \text{ cut off in the acceptance}$$

$$\frac{1}{f_{neutral}} \quad \begin{array}{l} \text{Correction for neutral hadrons included in the definition but not measured well:} \\ K_s^0, \Lambda, \bar{\Lambda}, K_L^0, n, \bar{n} \\ \text{Not trying to measure } K_s^0, \Lambda, \bar{\Lambda} \text{ in TPC – apply DCA cut to eliminate, correct for missing energy} \end{array}$$

$$f_{bg}^i(p_T) \quad \text{Correction for background not included in definition (e}^\mp\text{) or not measured easily event-by-event (} K_s^0, \Lambda, \bar{\Lambda} \text{)}$$

$$\frac{1}{f_{notID}} \quad \text{Correction for } \pi, K, p \text{ not identified}$$

$$eff(p_T^i) \quad \text{Correction for tracking efficiency}$$

$$E^{had} = \sqrt{p^2 + m^2} - m(\text{nucleons}) + \sqrt{p^2 + m^2} + m(\text{anti-nucleons}) + \sqrt{p^2 + m^2} (\text{all others})$$

Definition of energy to mimic the behavior of a calorimeter

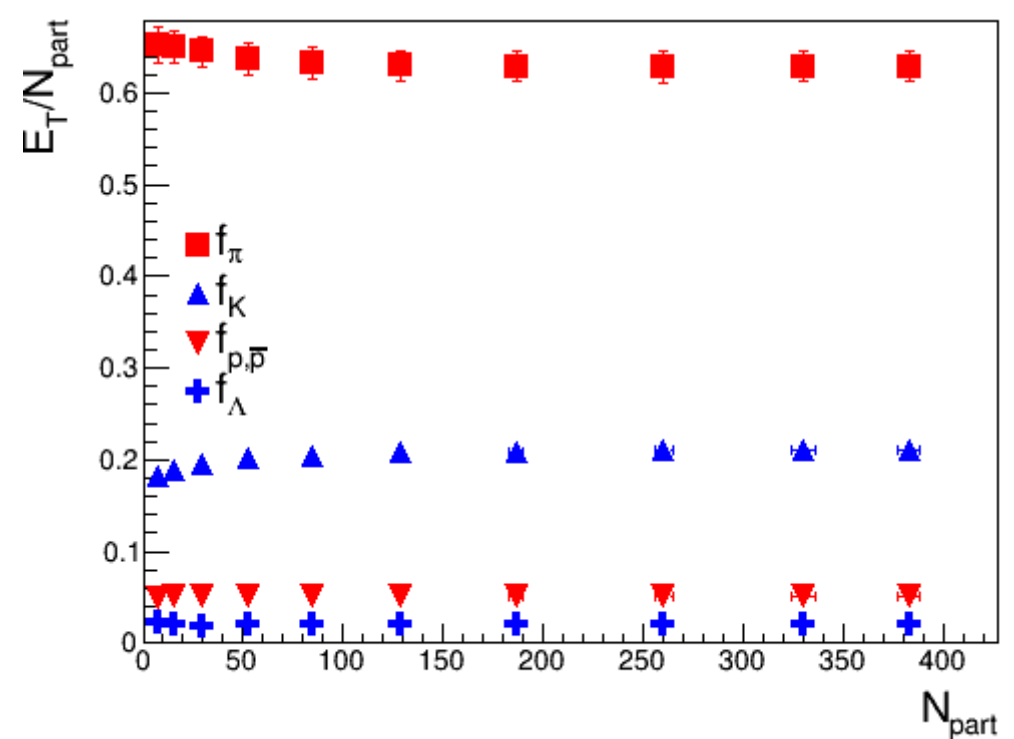
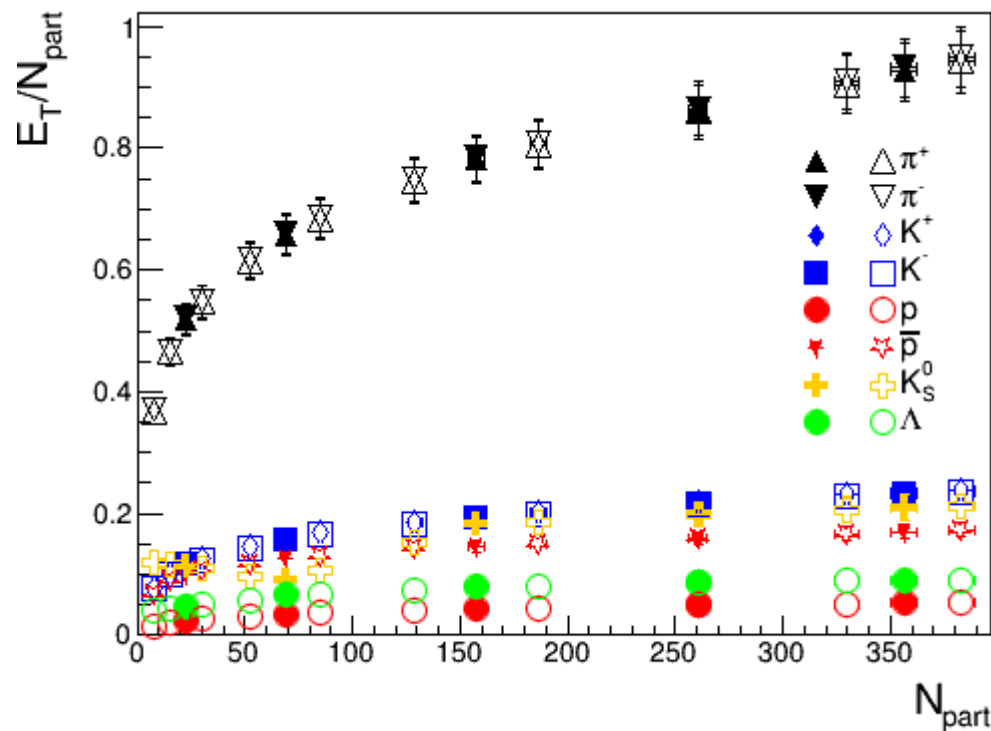
Calculation from spectra

- Use spectra data and use Blast wave fits to extrapolate to higher and lower p_T
- Three assumptions
$$E_T^n = E_T^p$$
$$E_T^{\bar{n}} = E_T^{\bar{p}}$$
$$K_L^0 = K_S^0$$
- Then, neglecting pseudorapidity dependence and assuming that the correction is the same for 900 GeV, 2.76 TeV, and 7 TeV:

$$E_T = E_T^{p, \bar{p}} + E_T^{n, \bar{n}} + E_T^K + E_T^\pi + E_T^{\Lambda, \bar{\Lambda}} + E_T^\eta$$

Everything else is negligible

Calculations from spectra

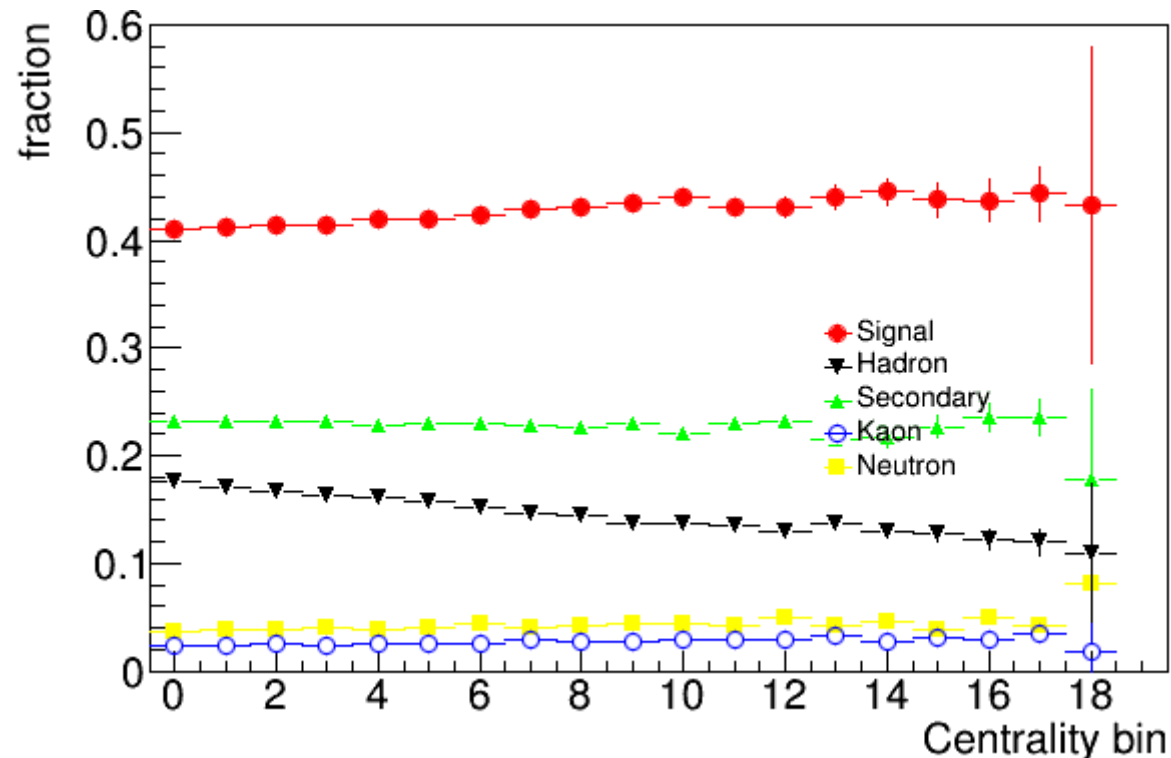


Neutral energy is not 1/3 of the energy. That's only true at low energies.

The calorimeter does not measure 1/3 of the energy. It only measures about 23%.

The distribution of energy is surprisingly centrality independent.

What does the EMCal measure?

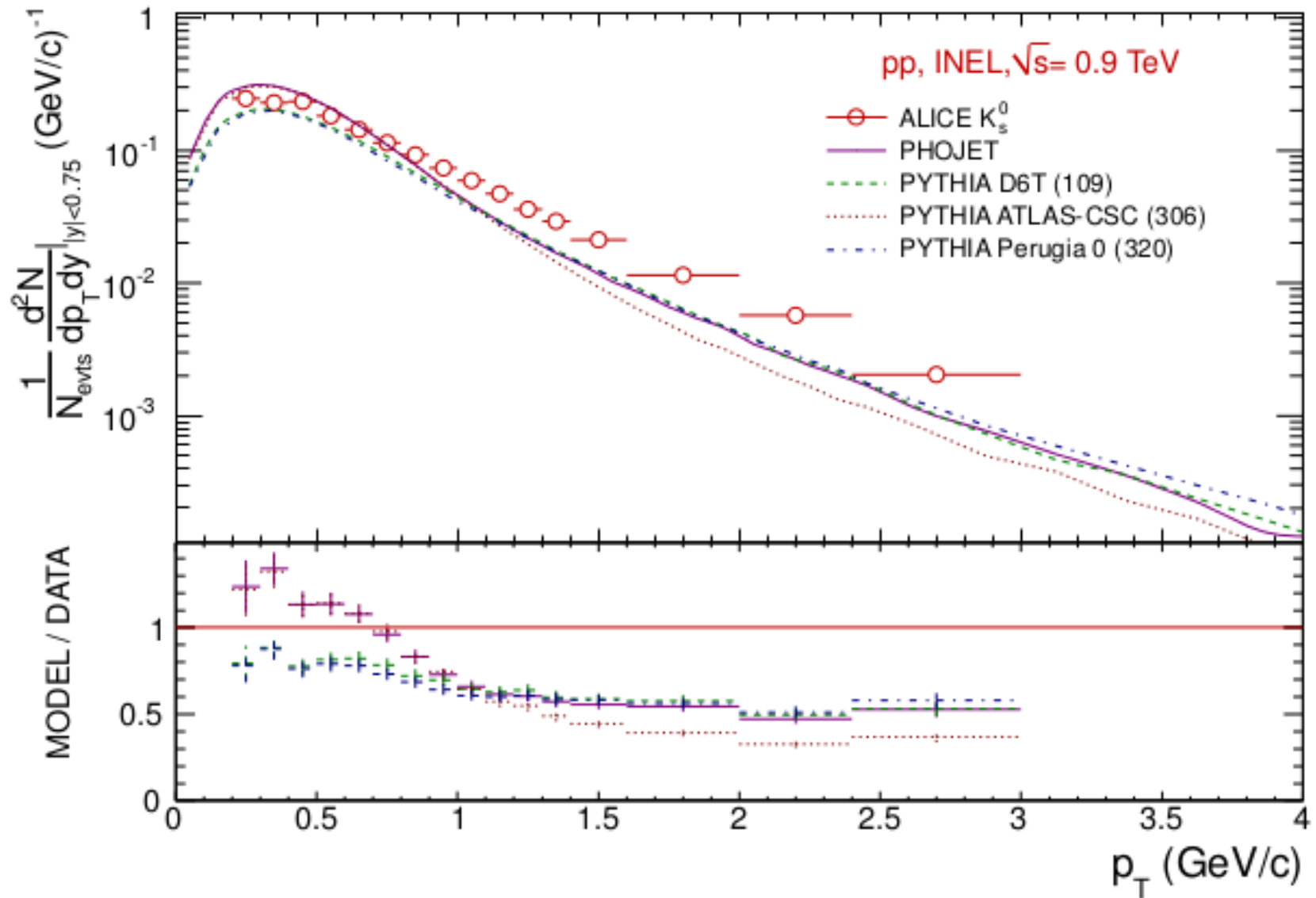


Note that this gets the fraction from kaons wrong. The fraction from kaons is actually about 10% of what we measure. Signal is actually ~30%.

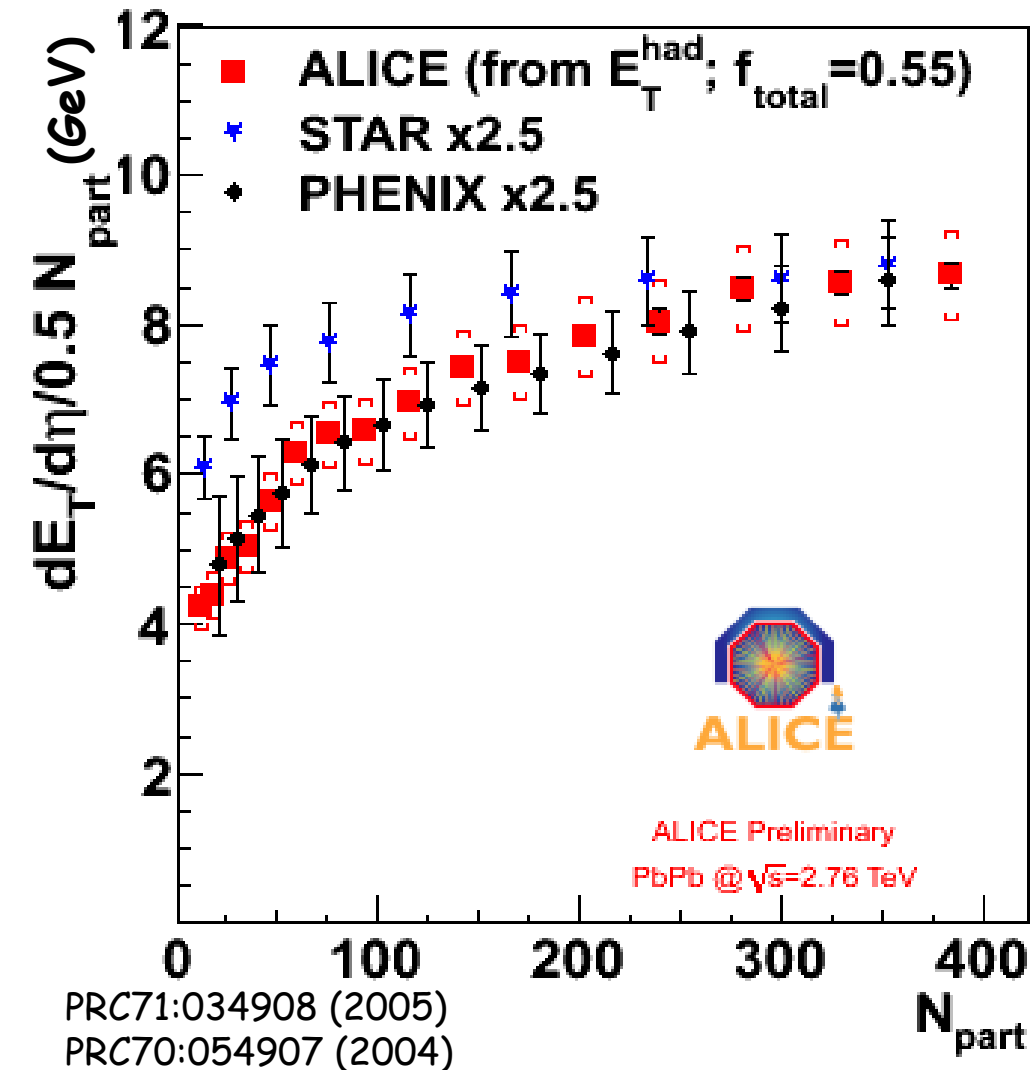
Kaon deposits

- There are several kaon decays into pi0's and pi0's decay mostly into photons
 - $K_S^0 \rightarrow \pi^0 \pi^0$ (30.7% B.R.)
 - $K^\pm \rightarrow \pi^\pm \pi^0$ (20.7% B.R.)
 - $K^\pm \rightarrow \pi^0 e^\pm \nu_e$ (5.1% B.R.)
 - $K^\pm \rightarrow \pi^0 \mu^\pm \nu_\mu$ (3.4% B.R.)
 - $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ (1.8% B.R.)
 - $K_L^0 \rightarrow \pi^0 \pi^0 \pi^0$ (19.5% B.R.)
 - $K_L^0 \rightarrow \pi^+ \pi^- \pi^0$ (12.5% B.R.).
- These will (mostly) not be matched to tracks
- Simulations are unreliable because of how far off simulations are for strange particles

Kaons – measured vs simulation



Transverse Energy



- E_{T}^{had} from charged hadrons directly measured by the tracking detectors
- f_{total} from MC to convert into total ET

- From RHIC to LHC

- ~2.5 increase in $dE_T/d\eta / (0.5 \cdot N_{part})$

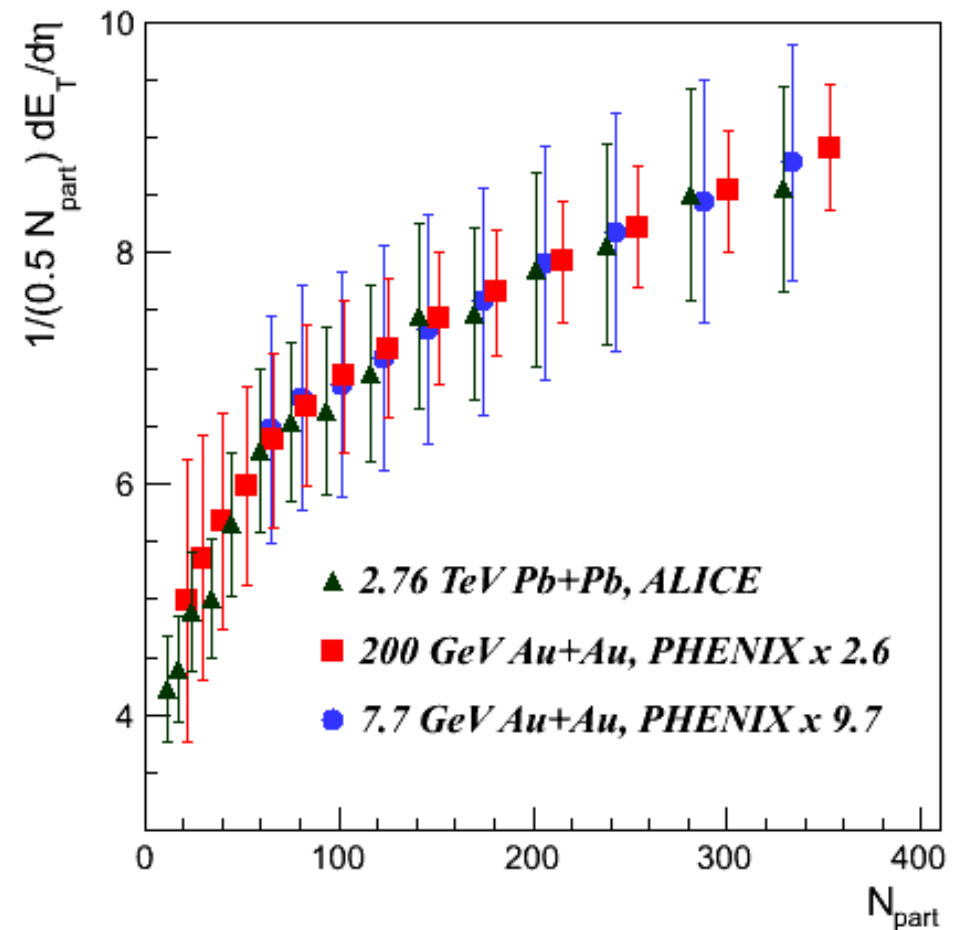
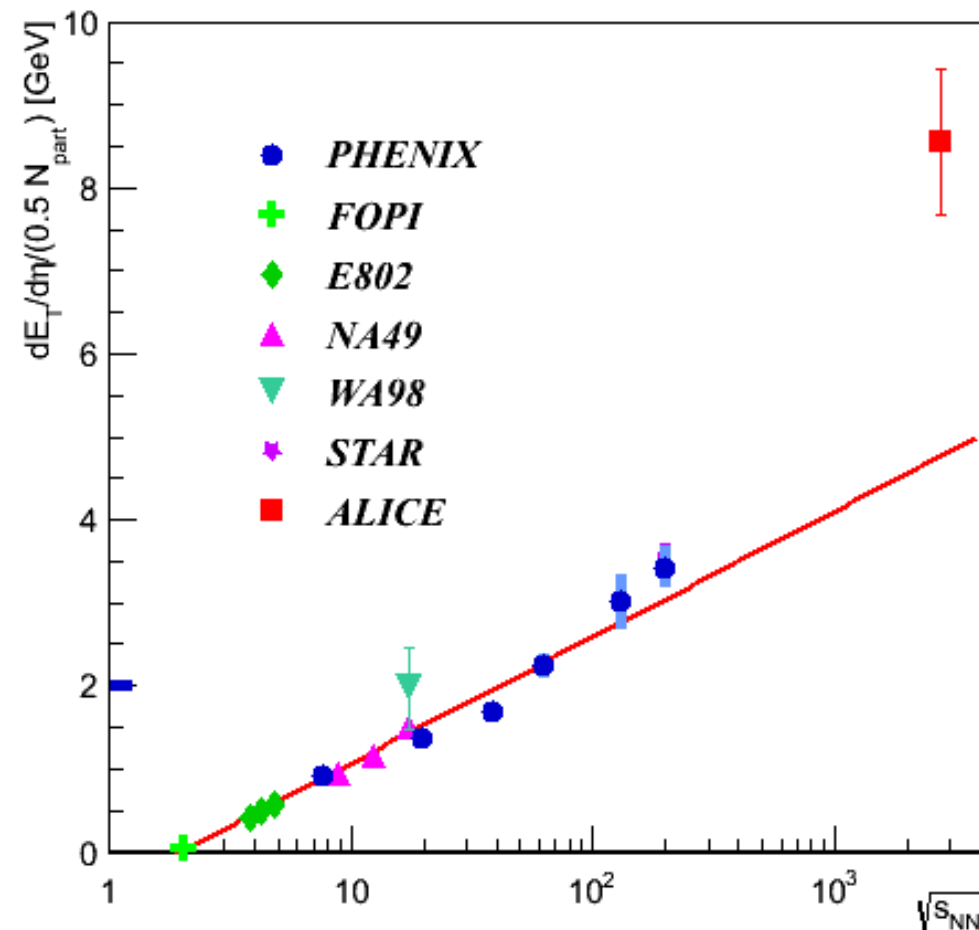
- Energy density (Bjorken)

$$\varepsilon = \frac{1}{\pi R^2 \tau} \frac{dE_t}{dy} \quad R = 1.12 A^{1/3} fm$$

- $\varepsilon \tau \sim 15 \text{ GeV}/(fm^2 c)$
RHIC: $\varepsilon \tau = 5.4 \pm 0.6 \text{ GeV}/(fm^2 c)$

Centrality dependence similar to
RHIC (PHENIX)

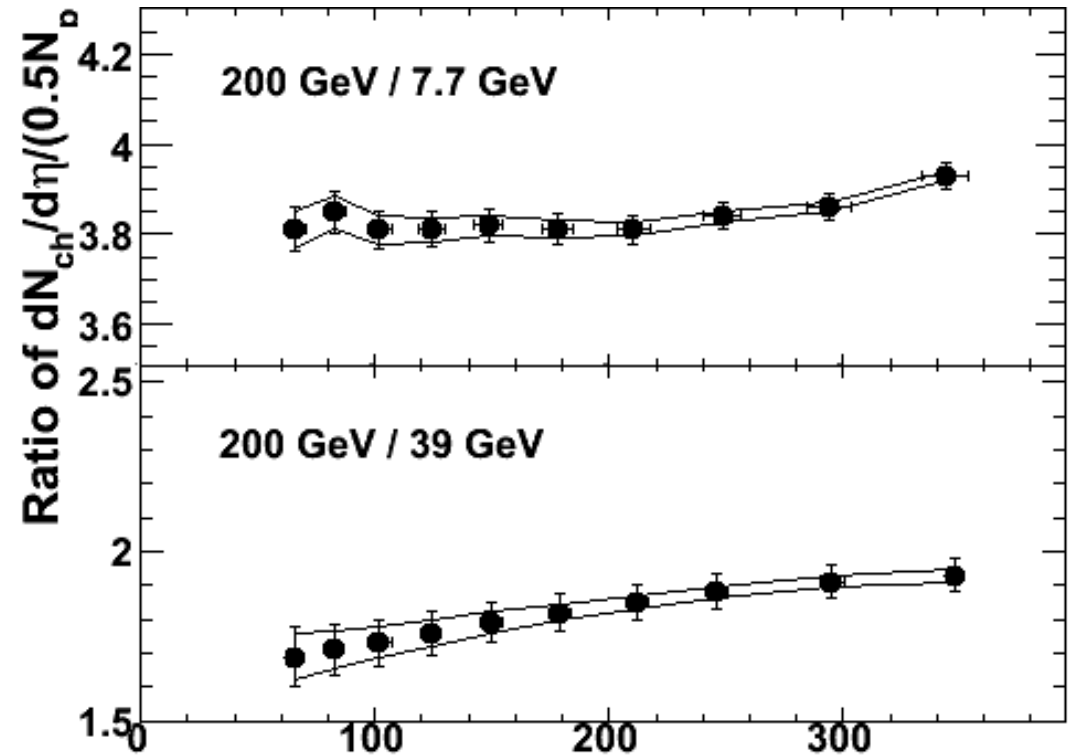
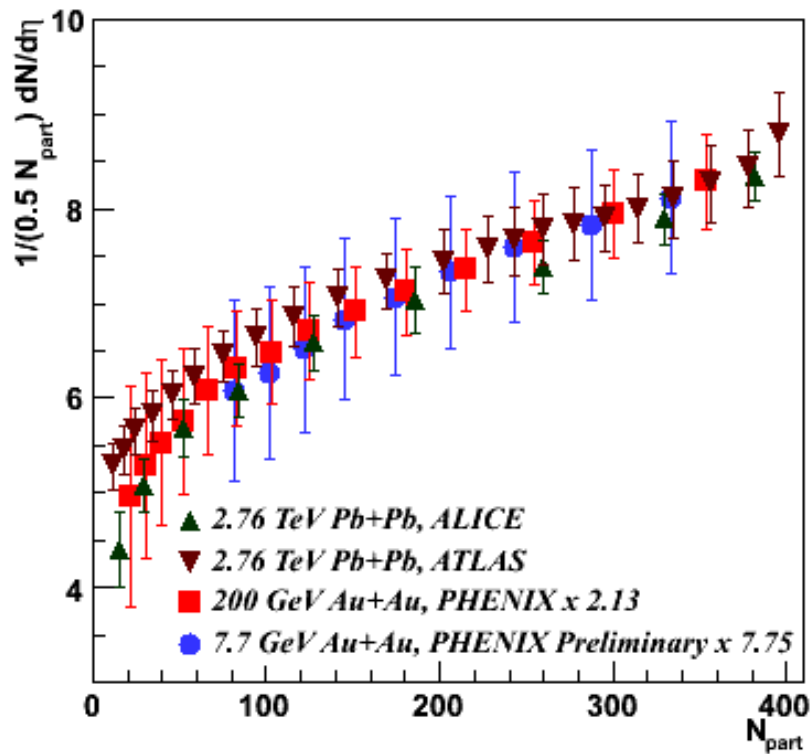
$dE_T/d\eta$



Also for transverse energy:

Approx. same centrality dependence at 7.7 GeV as at 2.76 TeV!

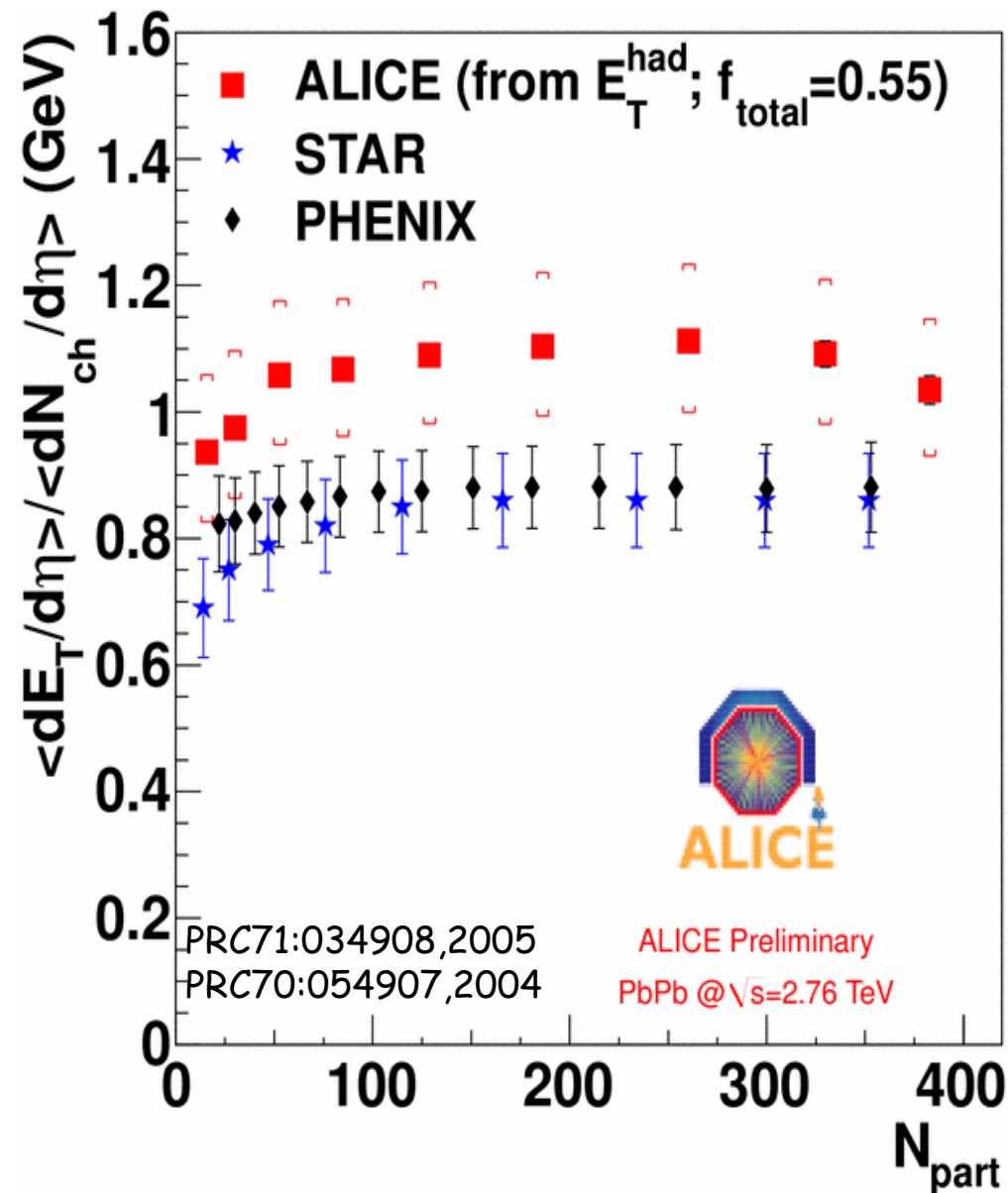
$$dN_{ch}/d\eta$$



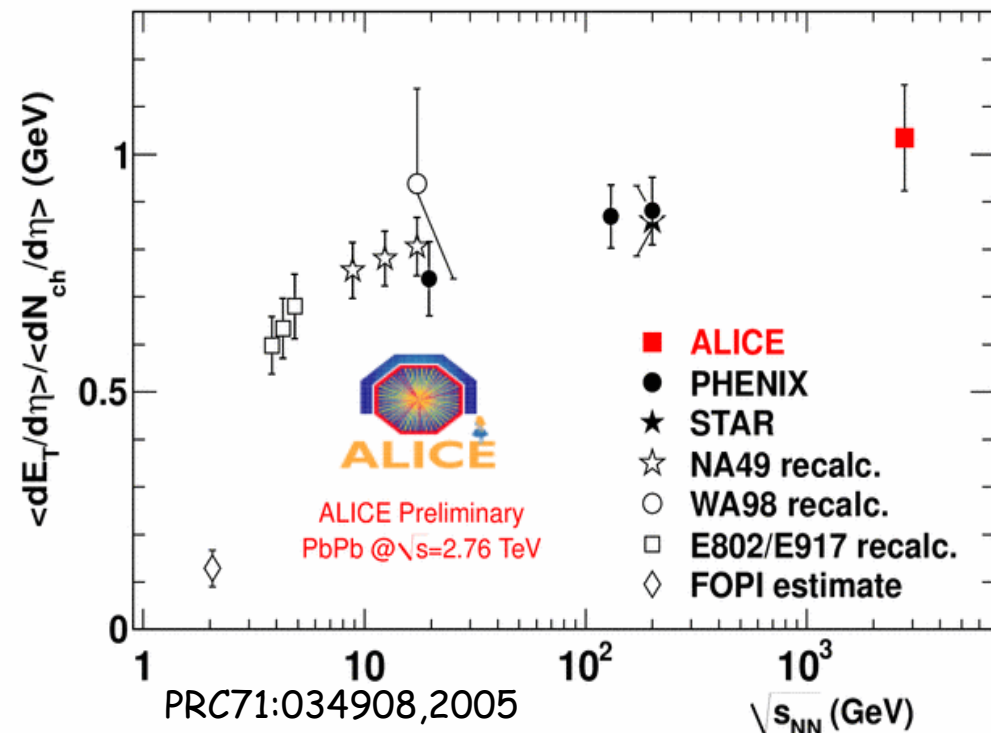
N.B.: Approx. same centrality dependence at 7.7 GeV as at 2.76 TeV!
 [note: no RHIC average here, just PHENIX..]

ET/Nch

- Consistent behavior for ET and Nch - Both increase with energy



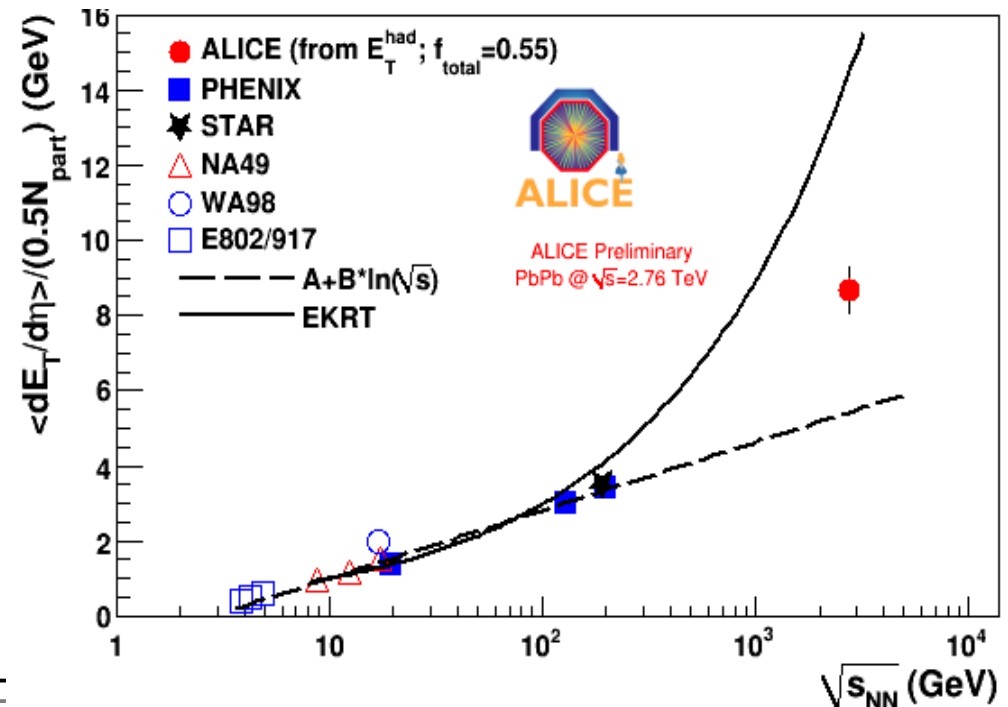
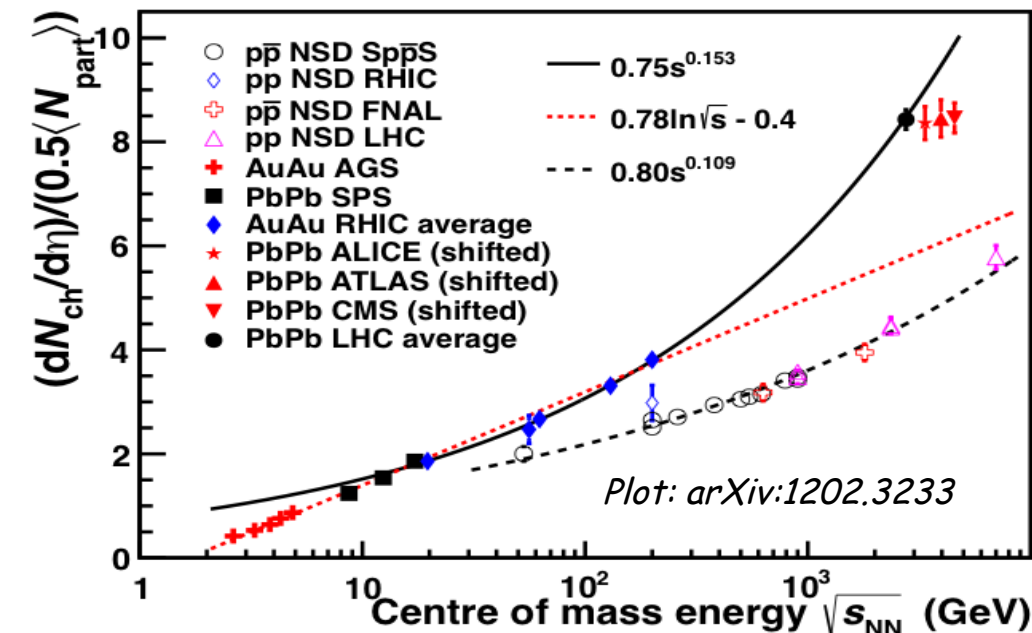
- Both show steady rise from peripheral to central
- ET/Nch \sim independent of centrality
- ET/Nch increases with energy



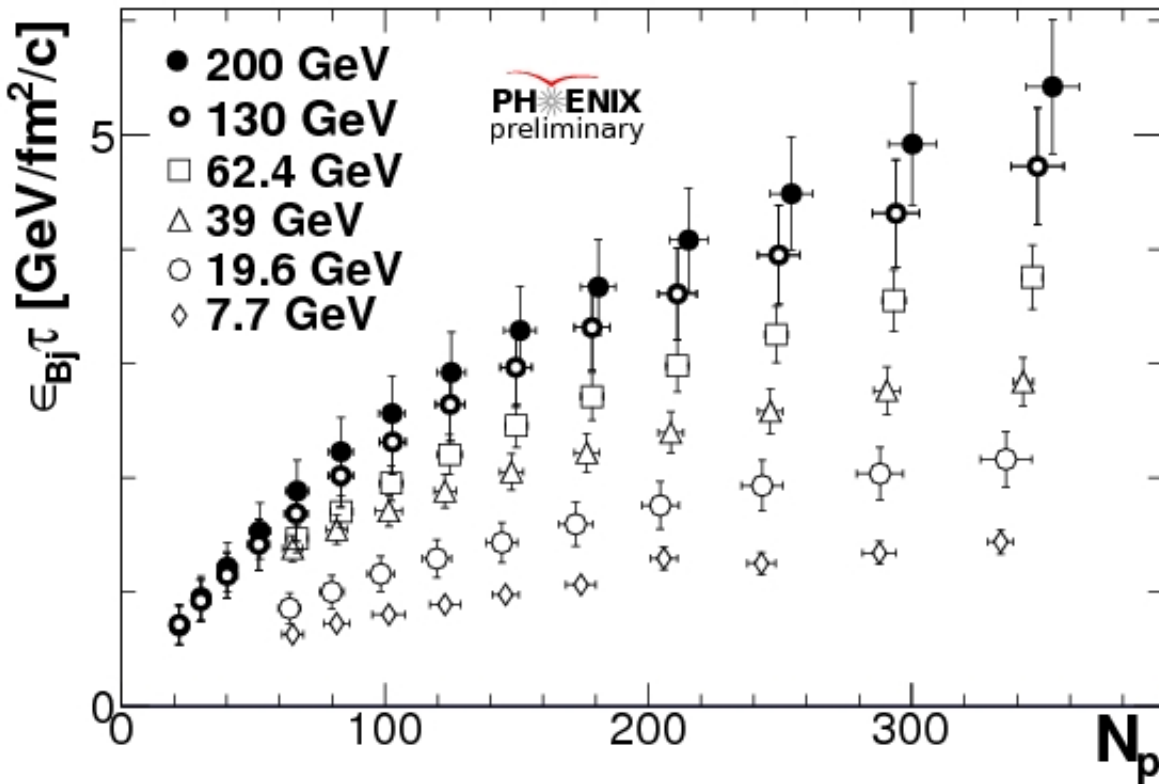
\sqrt{s} dependence: Nch & ET

- $dN_{ch}/d\eta/(0.5 \cdot N_{part}) \sim 8$
- $2.1 \times \text{RHIC}$
 $1.9 \times \text{pp (NSD) at 2.36 TeV}$
- growth with \sqrt{s} faster in AA than pp
- $dET/d\eta/(0.5 \cdot N_{part}) \sim 9$ in 0-5%
- Also increase of N_{part} (353 \rightarrow 383)
 $\rightarrow 2.7 \times \text{RHIC for } dET/d\eta$
(consistent with $\sim 20\%$ increase of $\langle pT \rangle$, and expectations from spectra)

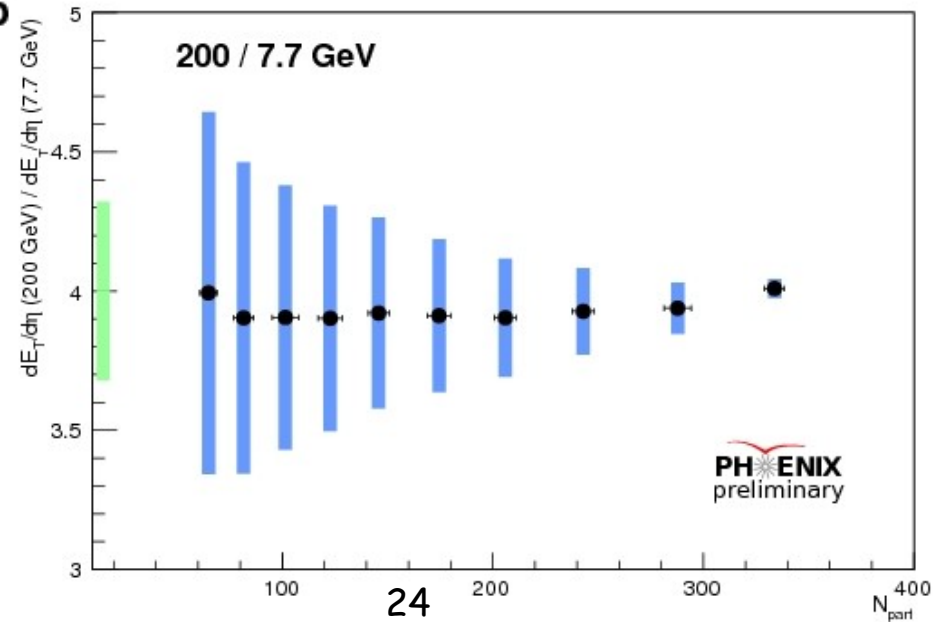
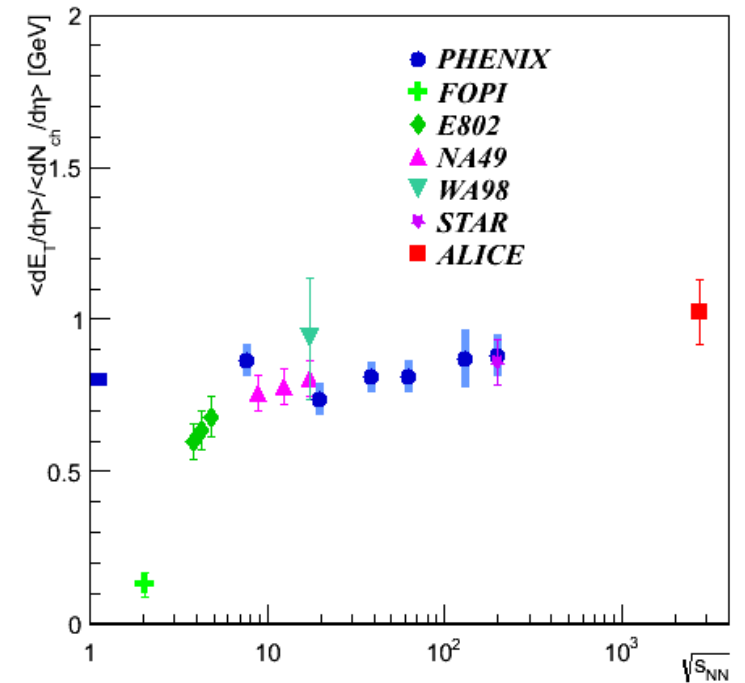
Grows with power of CM energy faster than simple logarithmic scaling extrapolated from lower energy



ET and Bjorken Energy Density



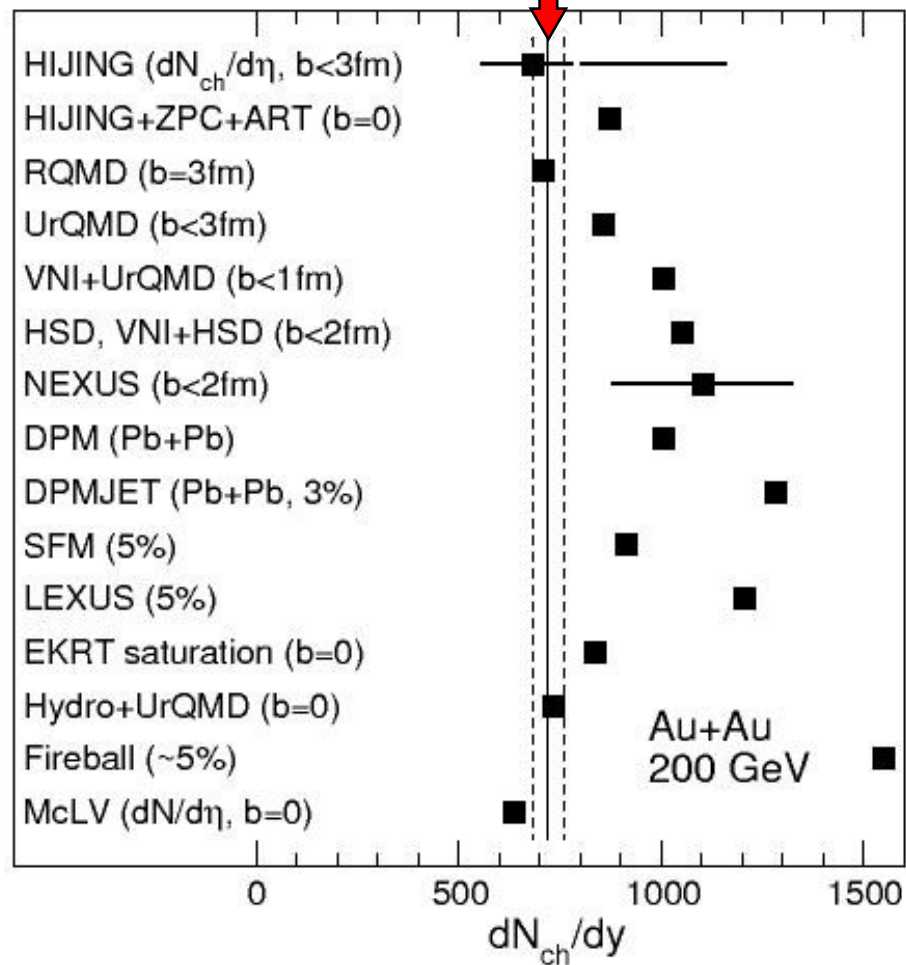
Smooth scaling throughout RHIC energy range



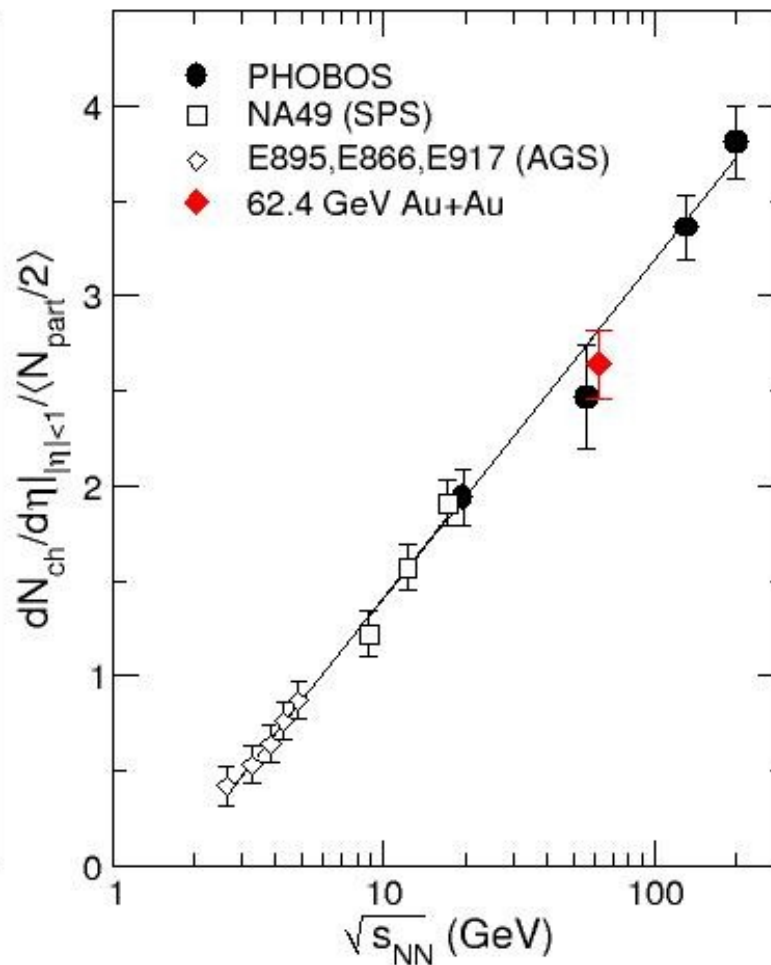
The End

From RHIC to LHC

Pre-RHIC theoretical predictions:

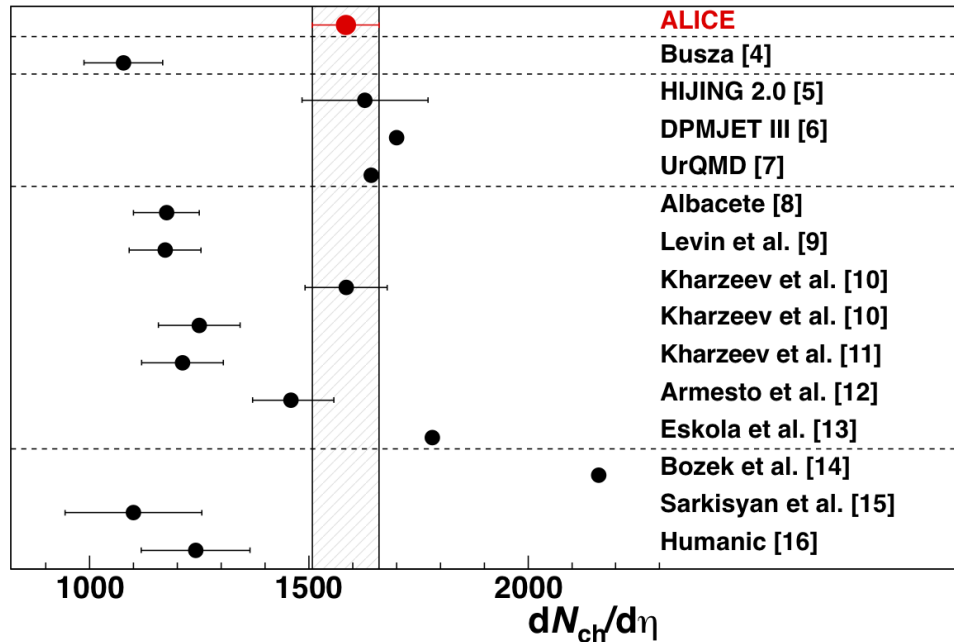


Seems straightforward to extrapolate to LHC, right..?



PHOBOS, Nucl. Phys. A747, 28 (2003)

First LHC HI Results: Charged Particle Multiplicity

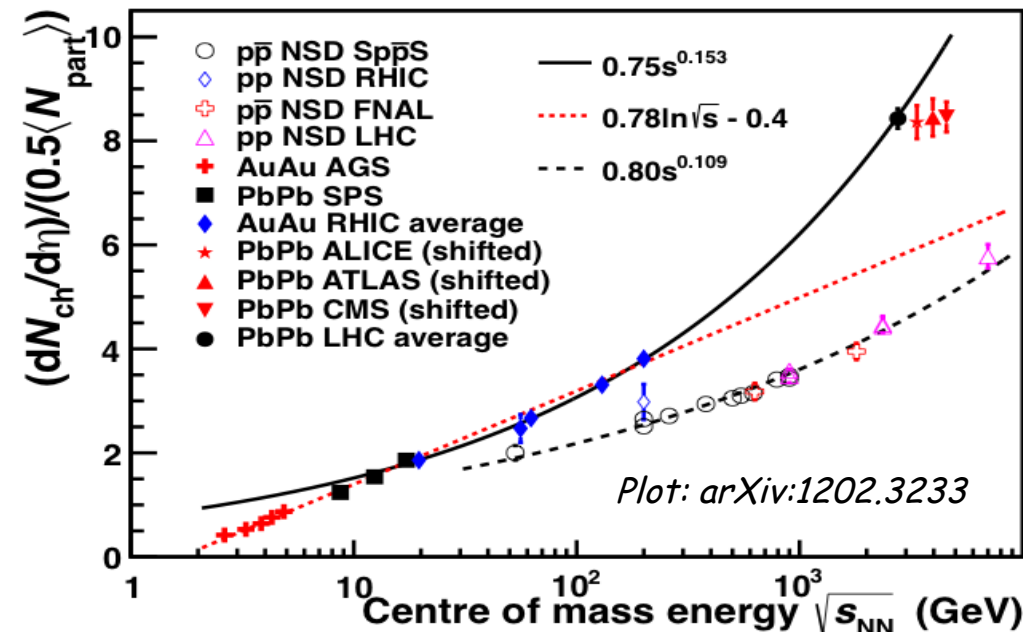


Excellent agreement also between LHC experiments!

Pb-Pb($\sqrt{s_{NN}} = 2.76$ TeV)
 $\rightarrow 1.9 \times$ p-p($\sqrt{s_{NN}} = 2.36$ TeV)
 \rightarrow nuclear amplification!
 $\rightarrow 2.1 \times$ RHIC (Au-Au($\sqrt{s_{NN}} = 0.2$ TeV))

5% most central events:
 $dN_{ch}/d\eta = 1584 \pm 4(\text{stat}) \pm 76(\text{sys})$

Predictions more spread around result (on high side of expectations) than at start of RHIC



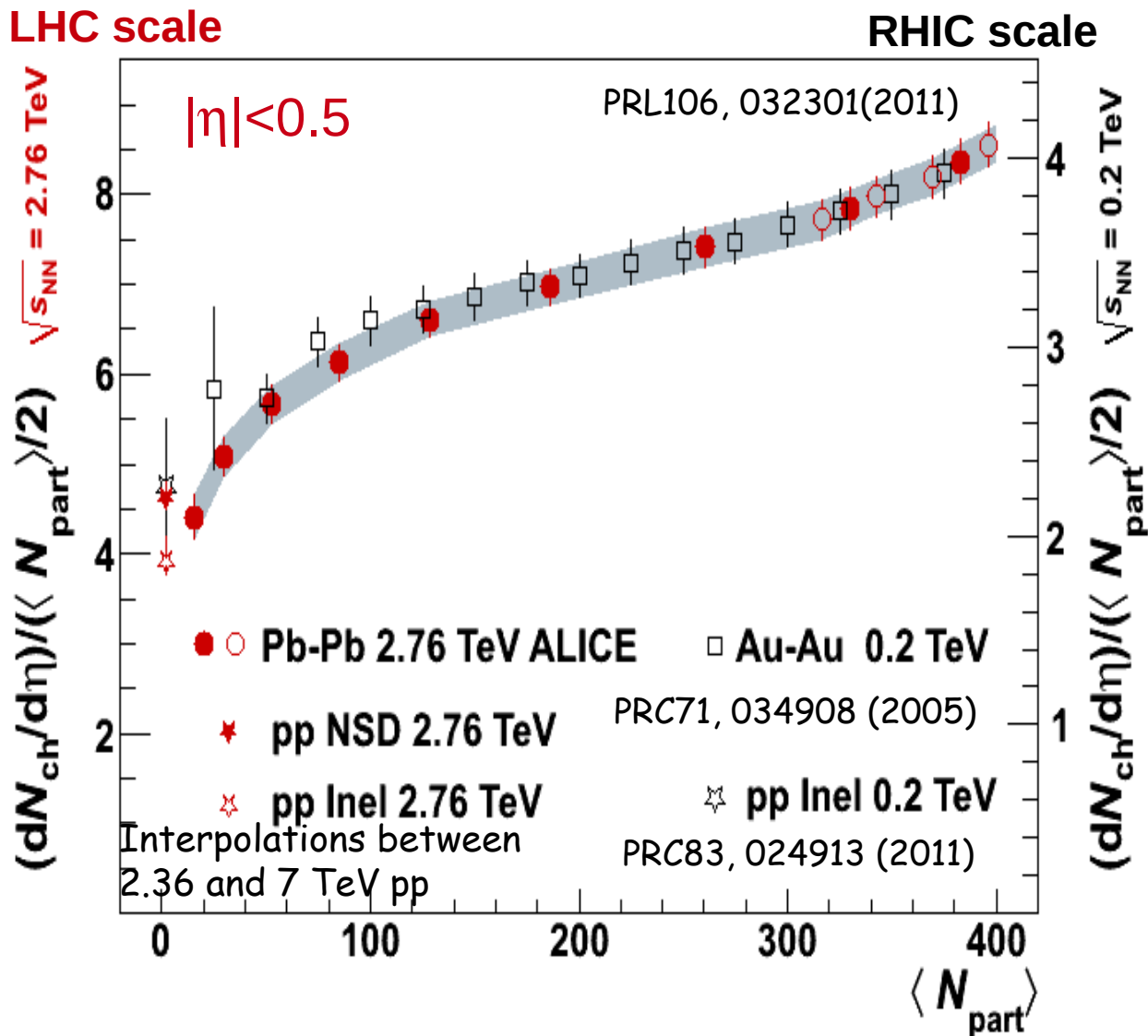
Summary

- The bulk properties of the system show a smooth transition throughout RHIC energy range and on to LHC energies
- LHC multiplicity (many predictions) and transverse energy (fewer comparisons) values higher than many predictions based on RHIC data
- Centrality dependence very similar from lowest RHIC to highest LHC \sqrt{s} (PHENIX & ALICE): “just” rapidity distribution narrowing/geometry?
- LHC Energy density $> 15 \text{ GeV/fm}^3 \rightarrow \sim 3\times \text{RHIC}$

Thanks to A. Milov, A. Toia, M. Floris, C. Nattrass, J. Mitchell for input/slides..

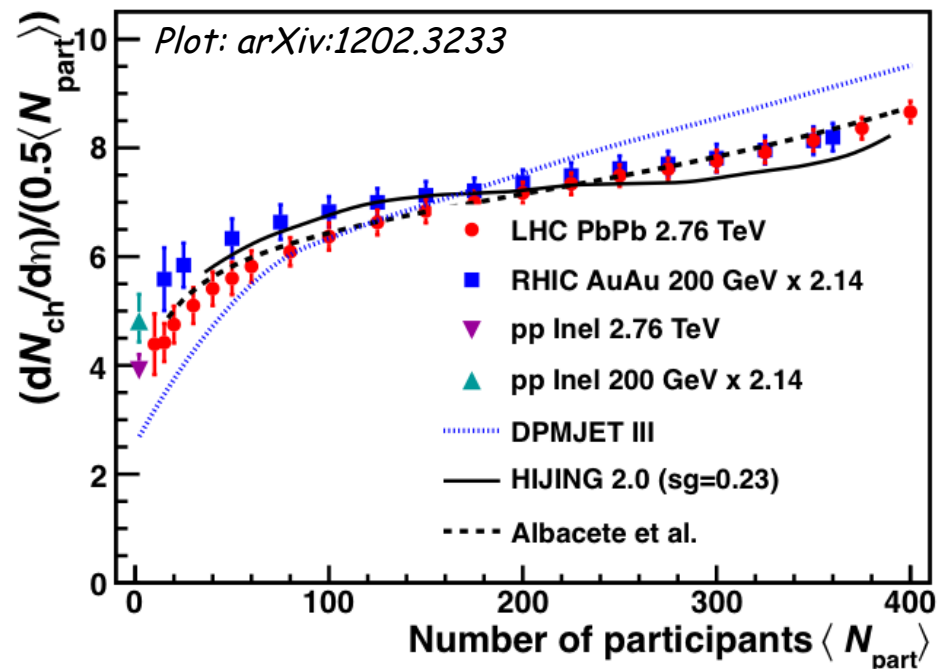
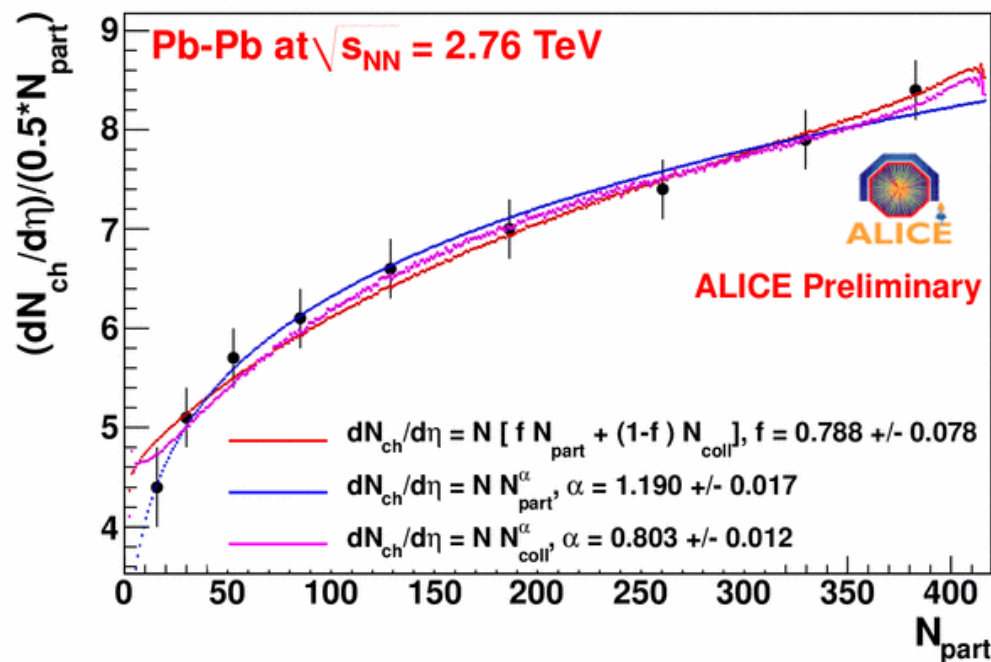
Multiplicity vs centrality

- Measurement based on tracklet reconstruction in SPD



- From RHIC to LHC
 - $dN_{ch}/d\eta \sim 1600$ for 0-5%: ~ 2.1 increase
 - Similar centrality dependence at 0.2 and 2.76 TeV for $N_{part} > 100$ (RHIC average)
 - Good "matching" to the pp point

Mid-rapidity $dN_{ch}/d\eta$ vs centrality



Multiplicity scaling with centrality:

- Stronger than N_{part}
- Different possible scalings (2 component, power laws) reproduce data
- Glauber fits not sensitive to choice of parameterization

Scaling similar to RHIC:

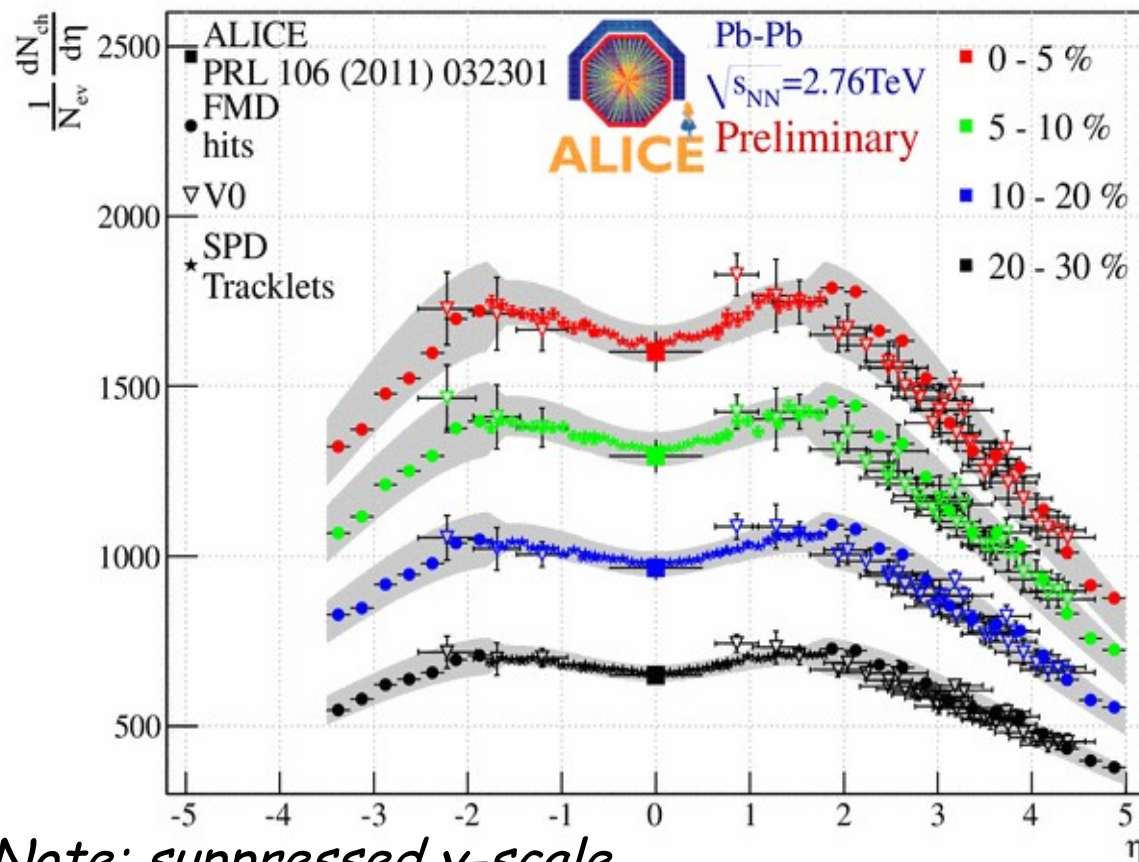
- Contribution of hard processes (N_{coll} scaling) the same as at lower energies..? Just geometry?

Observables at high rapidity

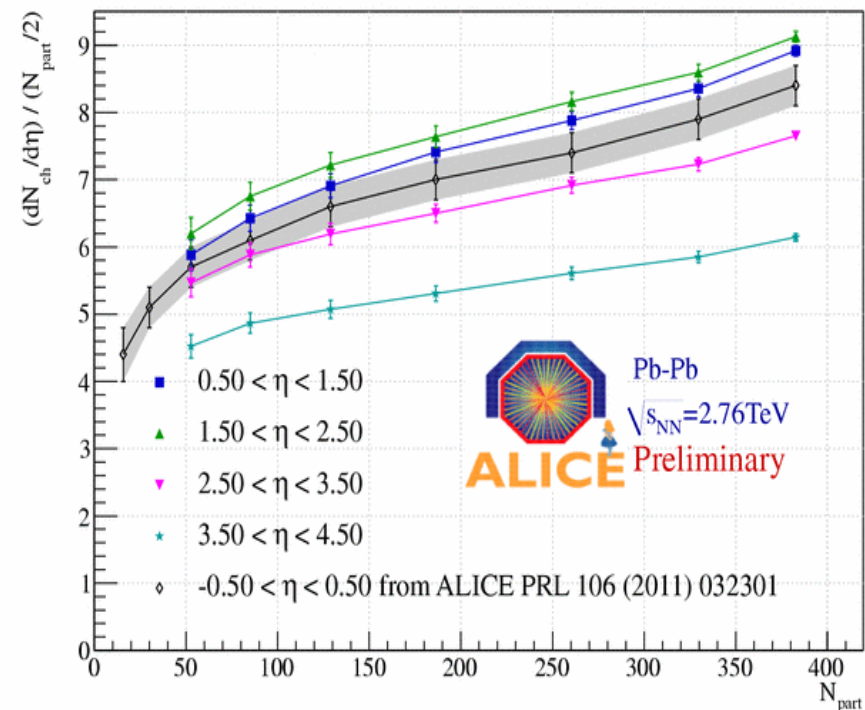
- $dN_{ch}/d\eta$ at forward rapidity
 - SPD: mid rapidity
 - VZERO, FMD

Phenomena at high rapidity

- properties of initial state (e.g. Color Glass condensate, gluon density, ...)
- energy and baryon stopping



similar trend for measured η bins



Note: suppressed y-scale..

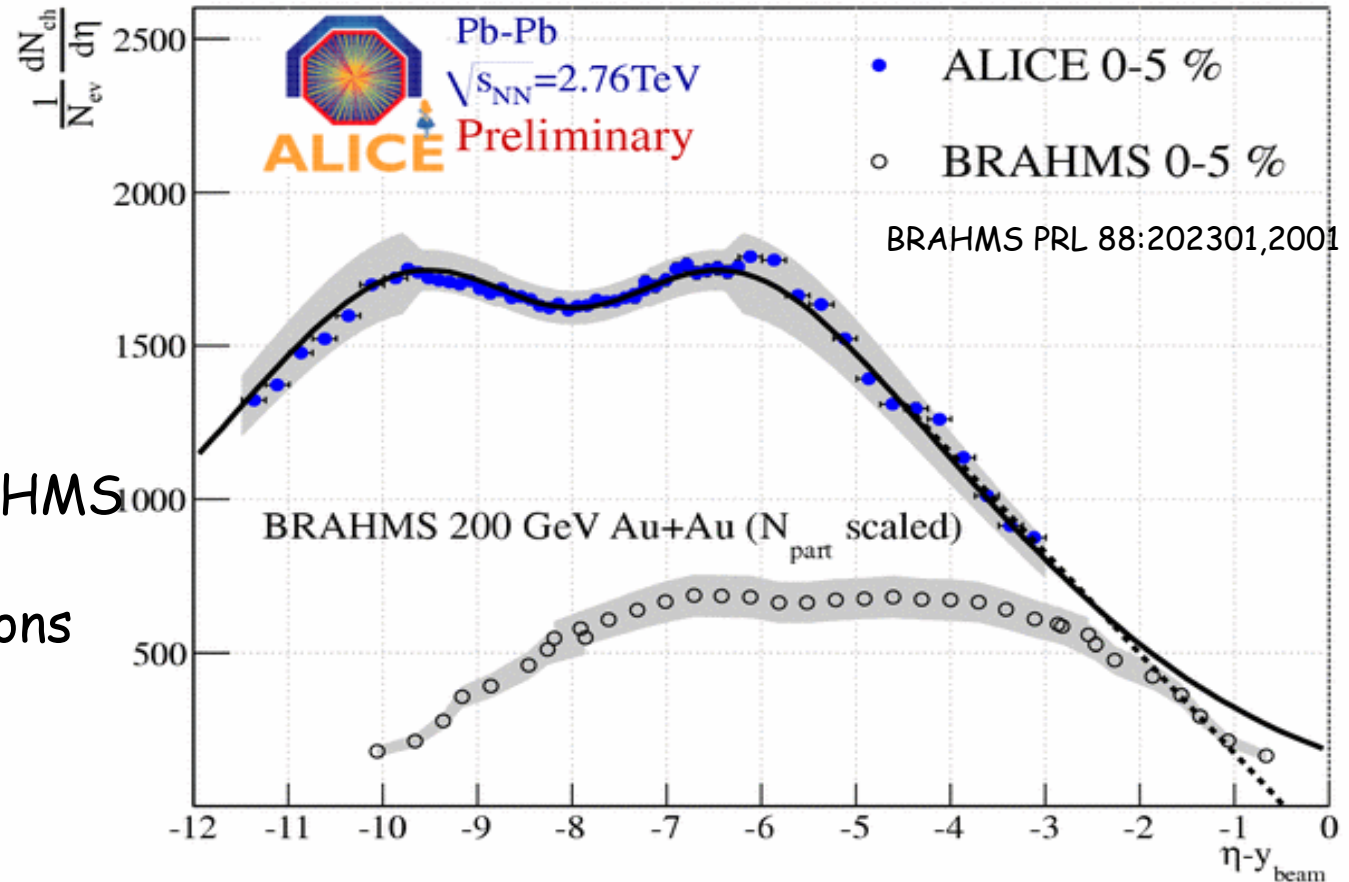
Christine Nattrass (U1K), Wayne State, Feb. 14, 2014

ALI-PREL-202

Extended longitudinal scaling

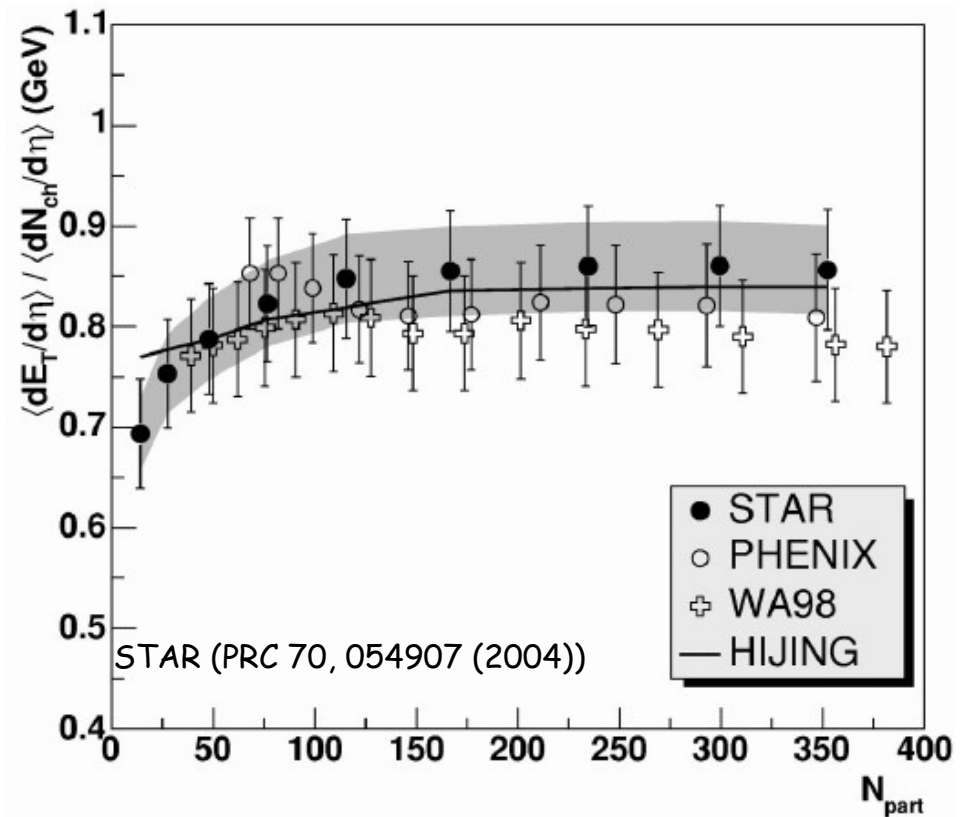
- Yields at high rapidity are energy-independent, when viewed in rest-frame of one of colliding nuclei
- longitudinal scaling could be present also for Pb+Pb at 2.76 TeV

Works also for $dN/d\eta$ because
 $y \approx \eta + \ln(pT/mT) \approx \eta$



ALI-PREL-194

ET Measurements

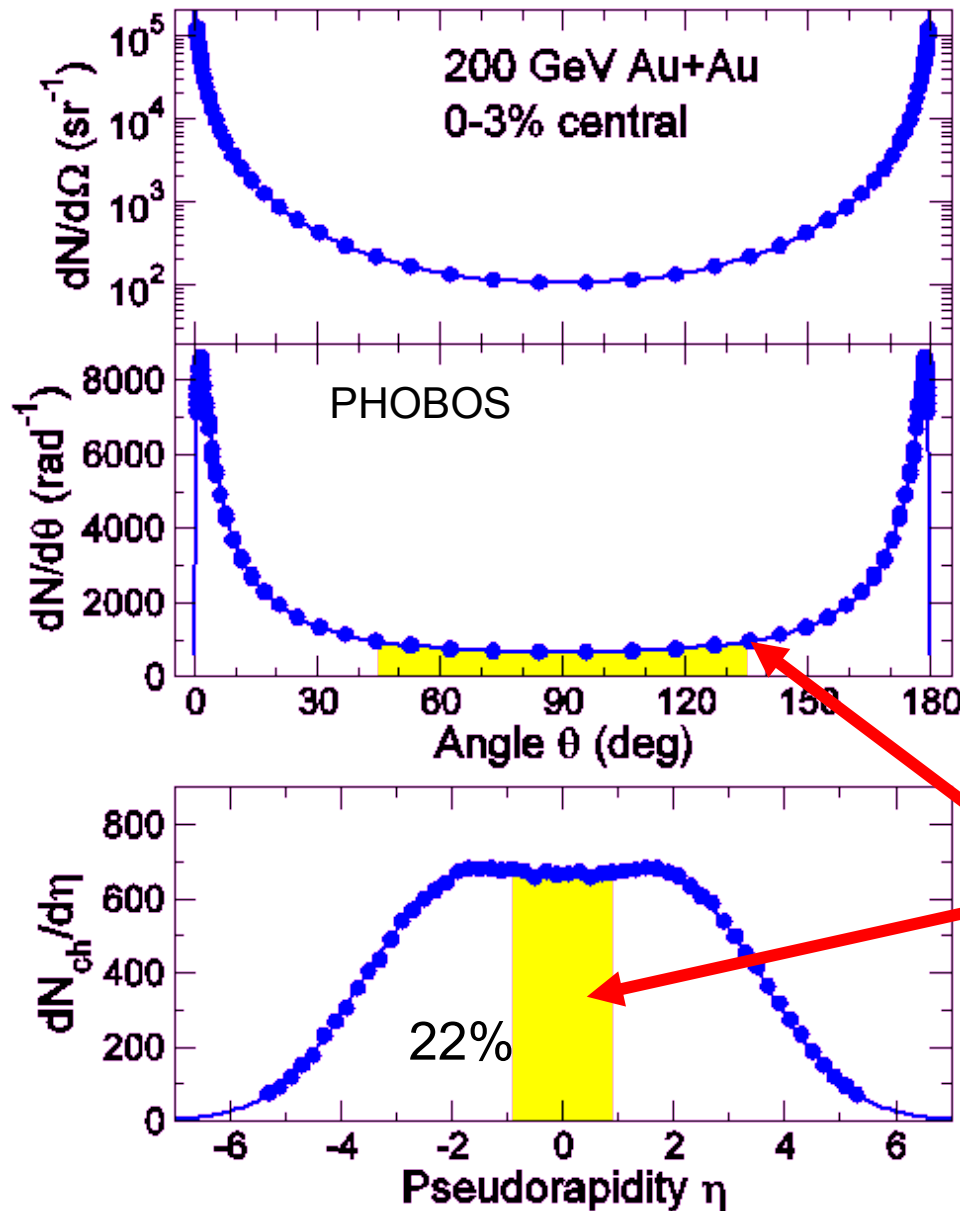


PHENIX PRC 71, 034908 (2005)

STAR: TPC + calorimetry
PHENIX: calorimetry only

- Centrality shape scales with incident beam energy
- Steady rise from peripheral to central a la Nch

Where do the particles go?



$$y = 0.5 \ln(E+pz)/(E-pz)$$

$$y_{\text{beam}} = \ln(\sqrt{s}/m_p)$$

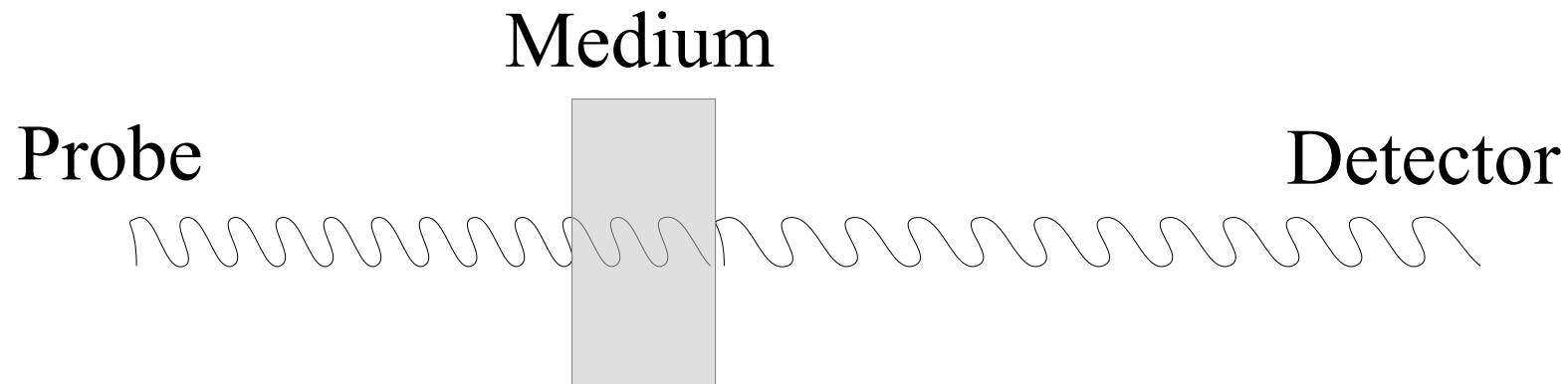
$$\sinh(\eta) = mT/pT \sinh(y)$$

Only ~22% of all emitted particles have $p_T > p_L$

Measurements at mid-rapidity carry information about the most dense region in the collision

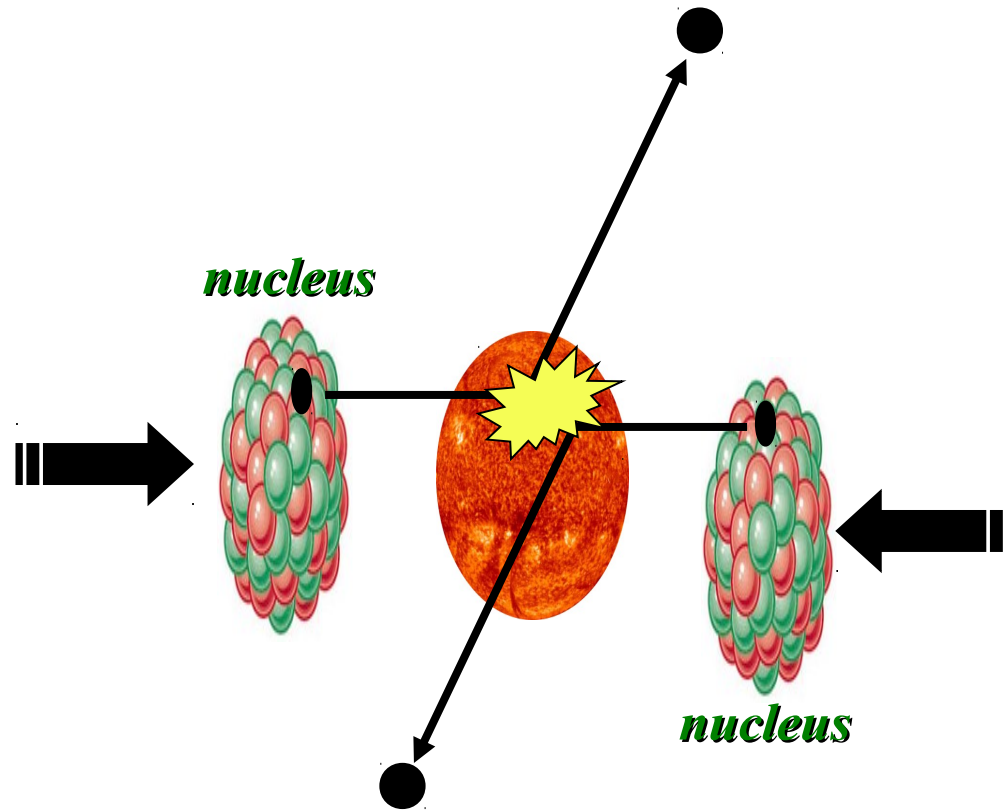
=> Let's focus on this region next..

Probing the Quark Gluon Plasma



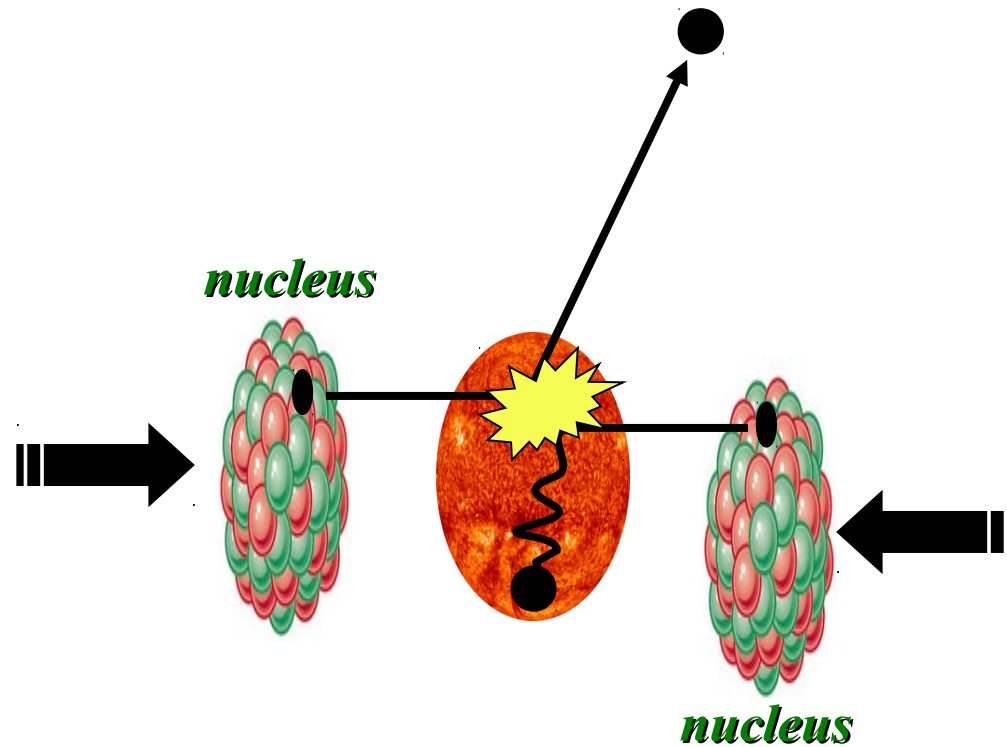
Want a probe which traveled through the collision
QGP is very short-lived ($\sim 1-10$ fm/c) \rightarrow
cannot use an external probe

Probes of the Quark Gluon Plasma



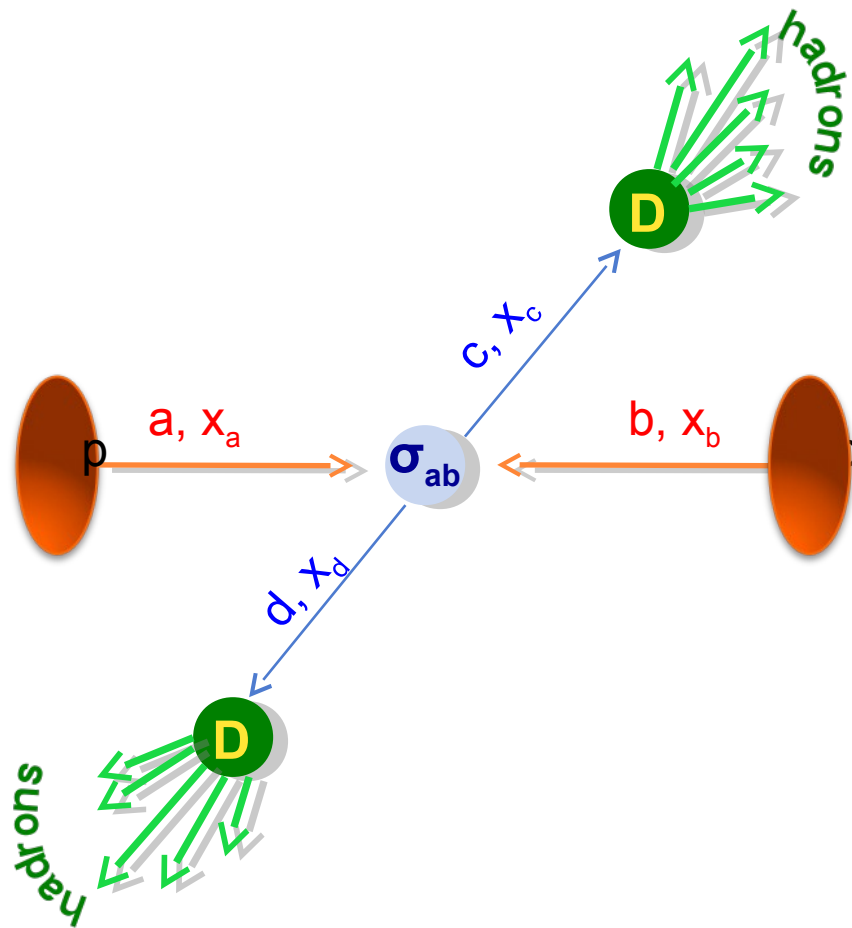
Want a probe which traveled through the medium
QGP is short lived \rightarrow need a probe created in the collision

Probes of the Quark Gluon Plasma

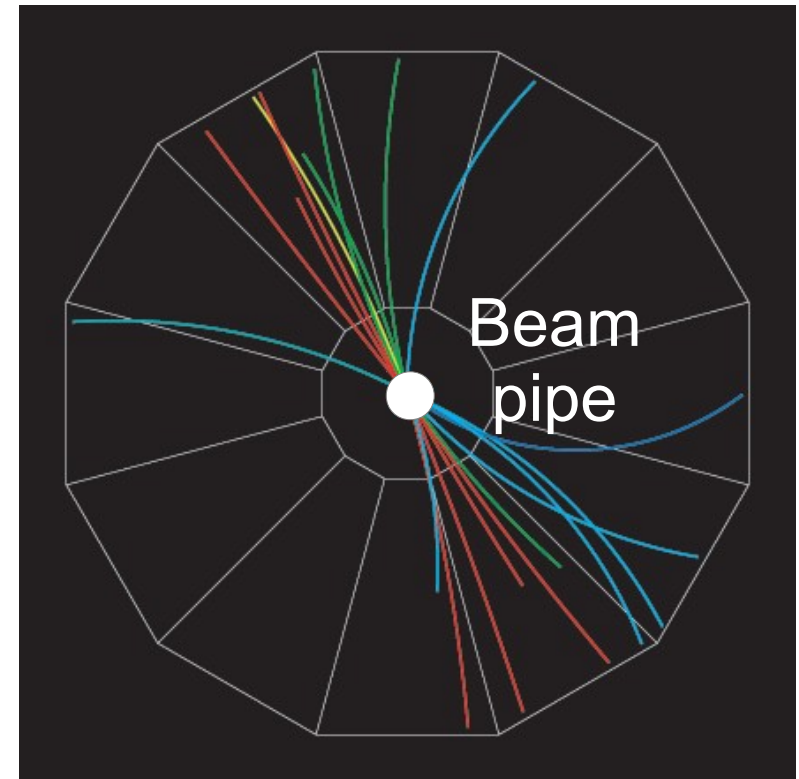


Want a probe which traveled through the medium
QGP is short lived \rightarrow need a probe created in the collision
We expect the medium to be dense \rightarrow absorb/modify probe

Jets

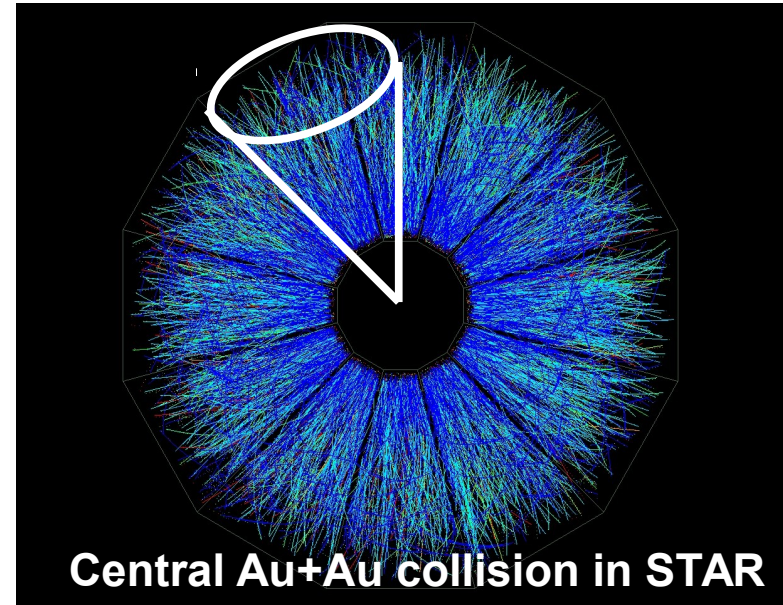
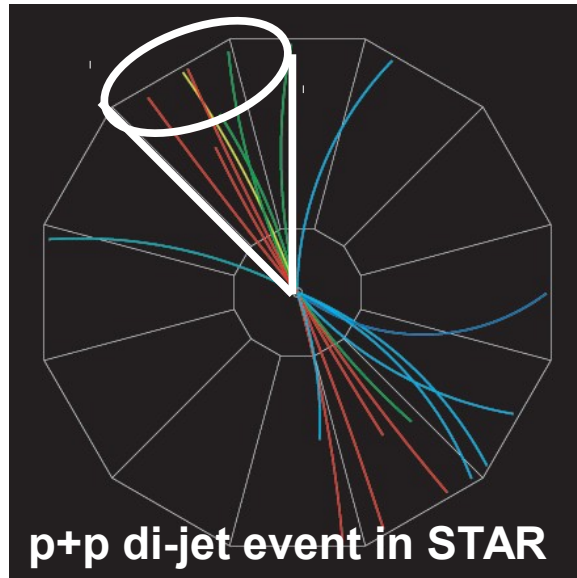


p+p → dijet



Jets – hard parton scattering leads to back-to-back quarks or gluons, which then fragment as a columnated spray of particles

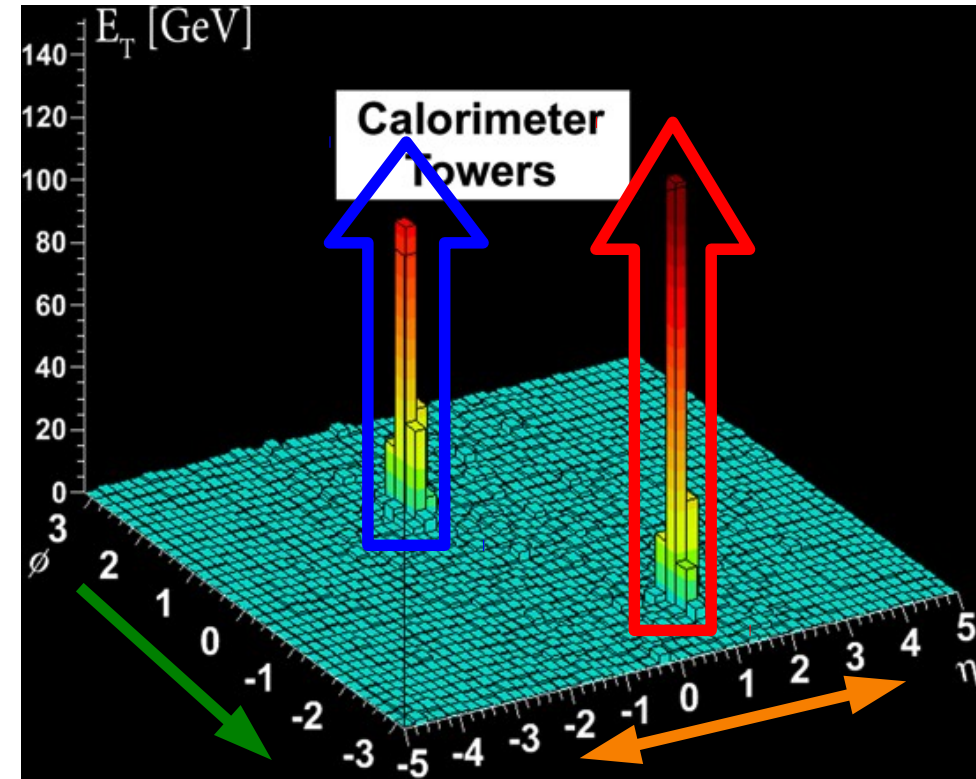
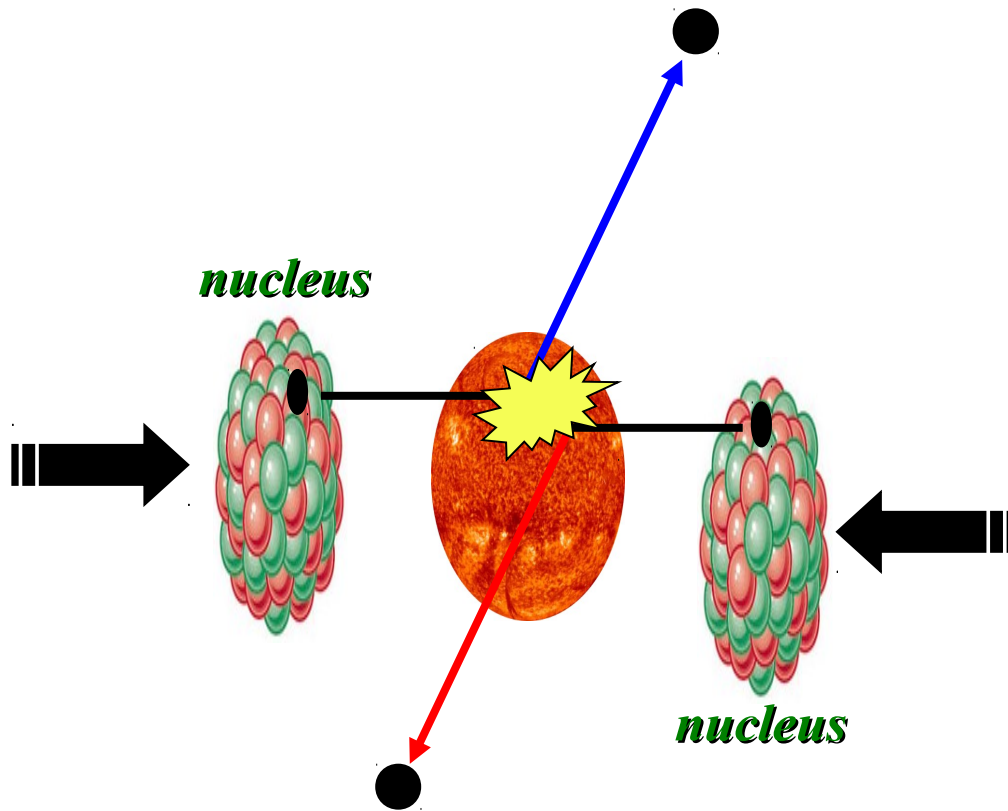
Jet reconstruction



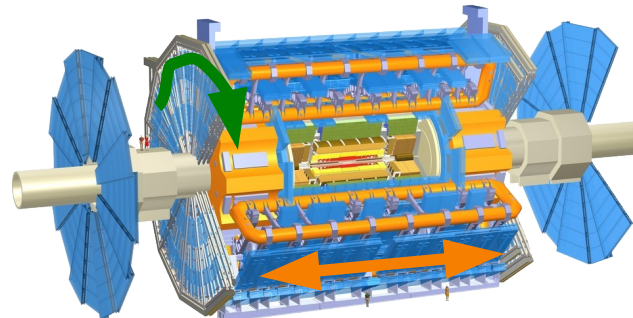
- Identify all of the particles in the jet \rightarrow parton energy, momentum
- Difficult in heavy ion collisions – but possible!

Jets

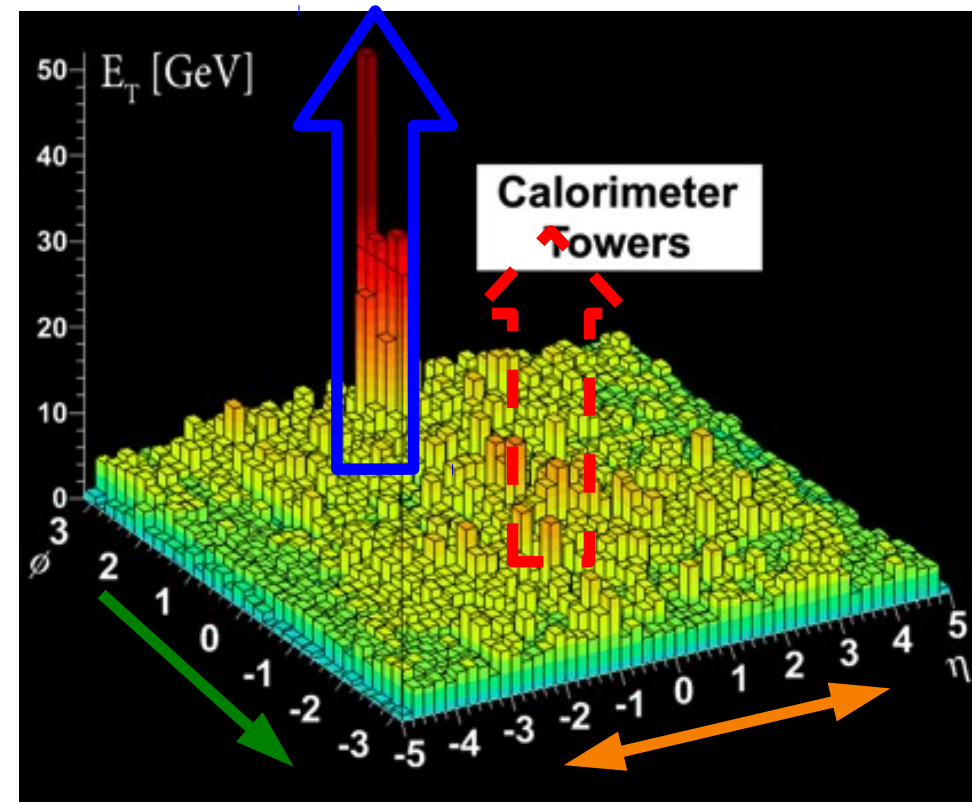
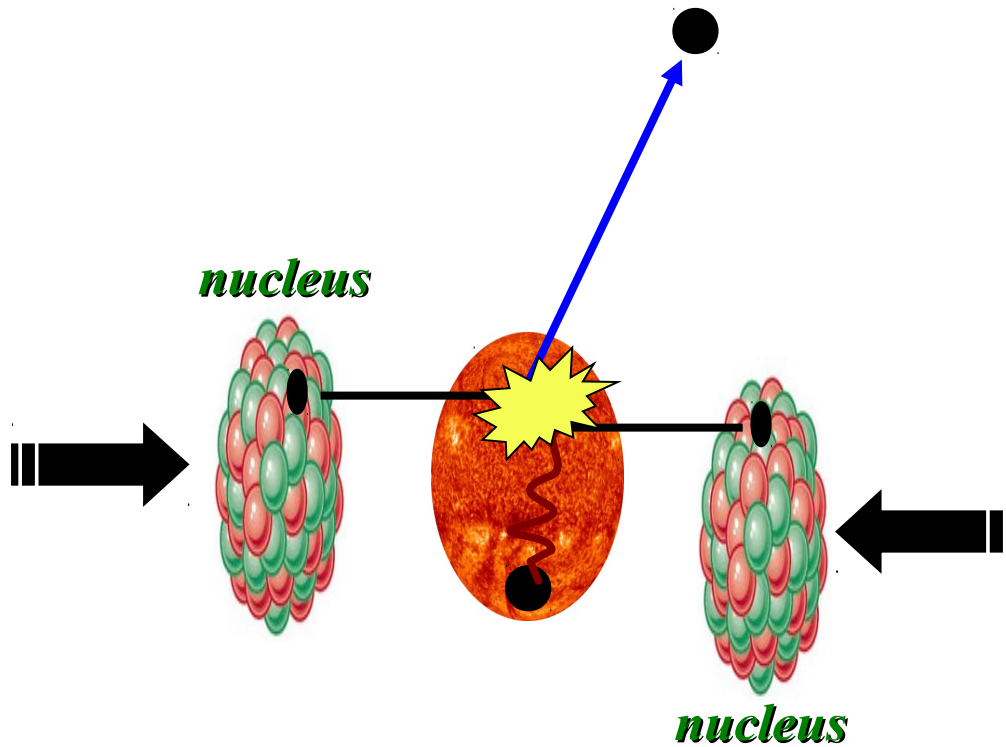
Phys.Rev.Lett. 105 (2010) 252303



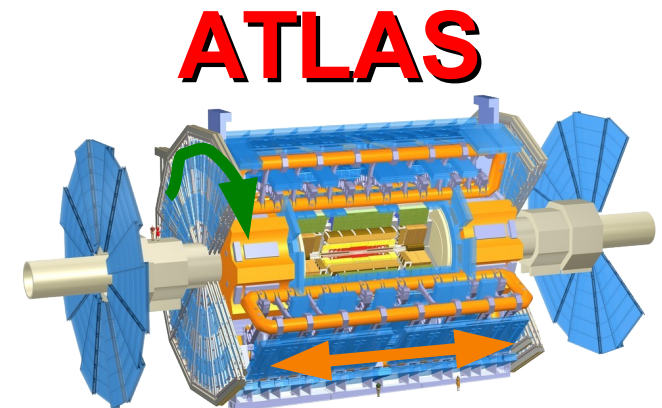
ATLAS



Quenched jets

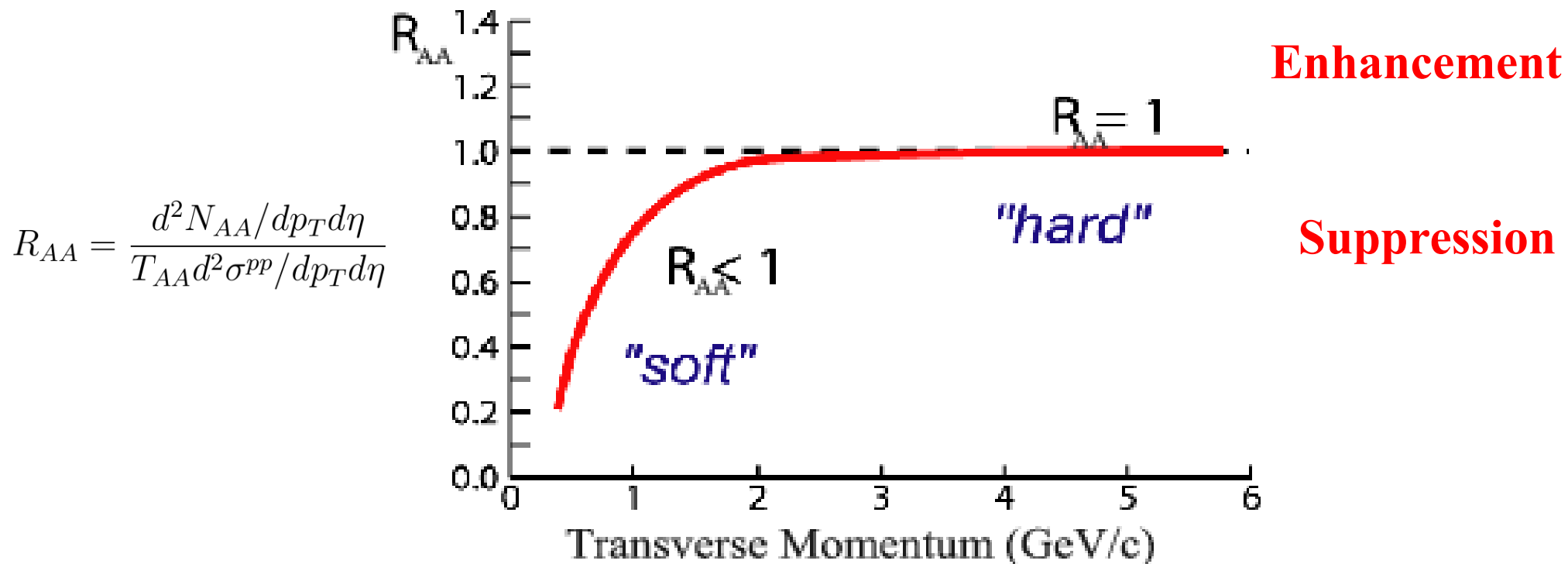


- One of the jets is absorbed by the medium
- The quark or gluon has equilibrated with the medium
- Phys. Rev. Lett. 105, 252303 (2010)



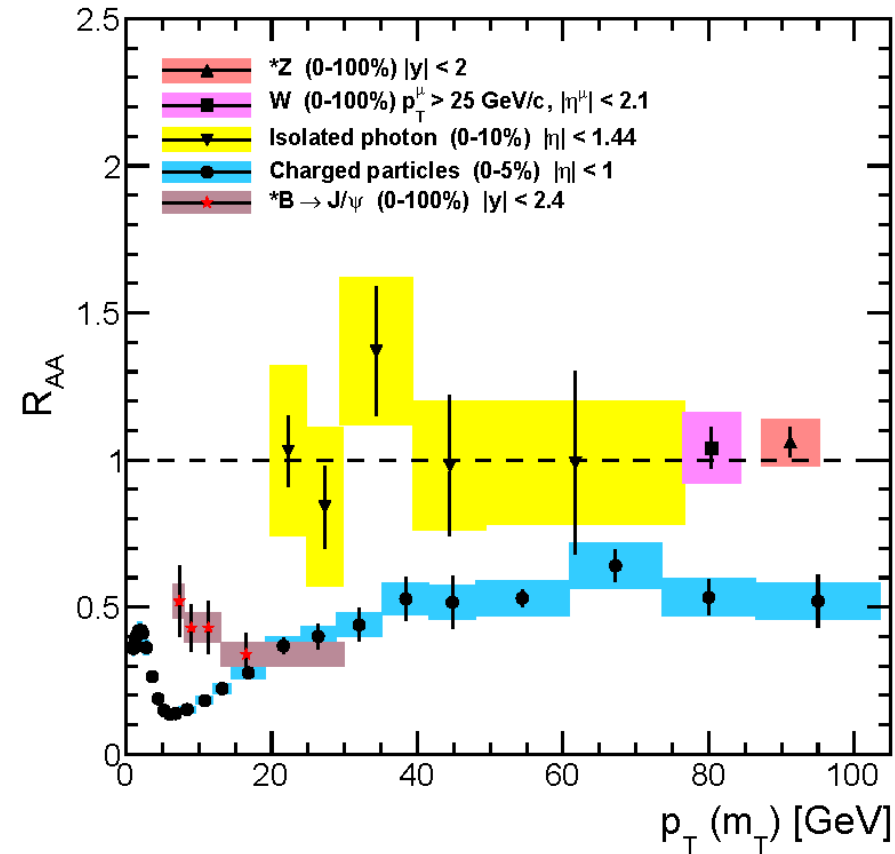
Nuclear modification factor

- Measure spectra of probe (jets) and compare to those in p+p collisions or peripheral A+A collisions
- If high- p_T probes (jets) are suppressed, this is evidence of jet quenching



Nuclear modification factor R_{AA}

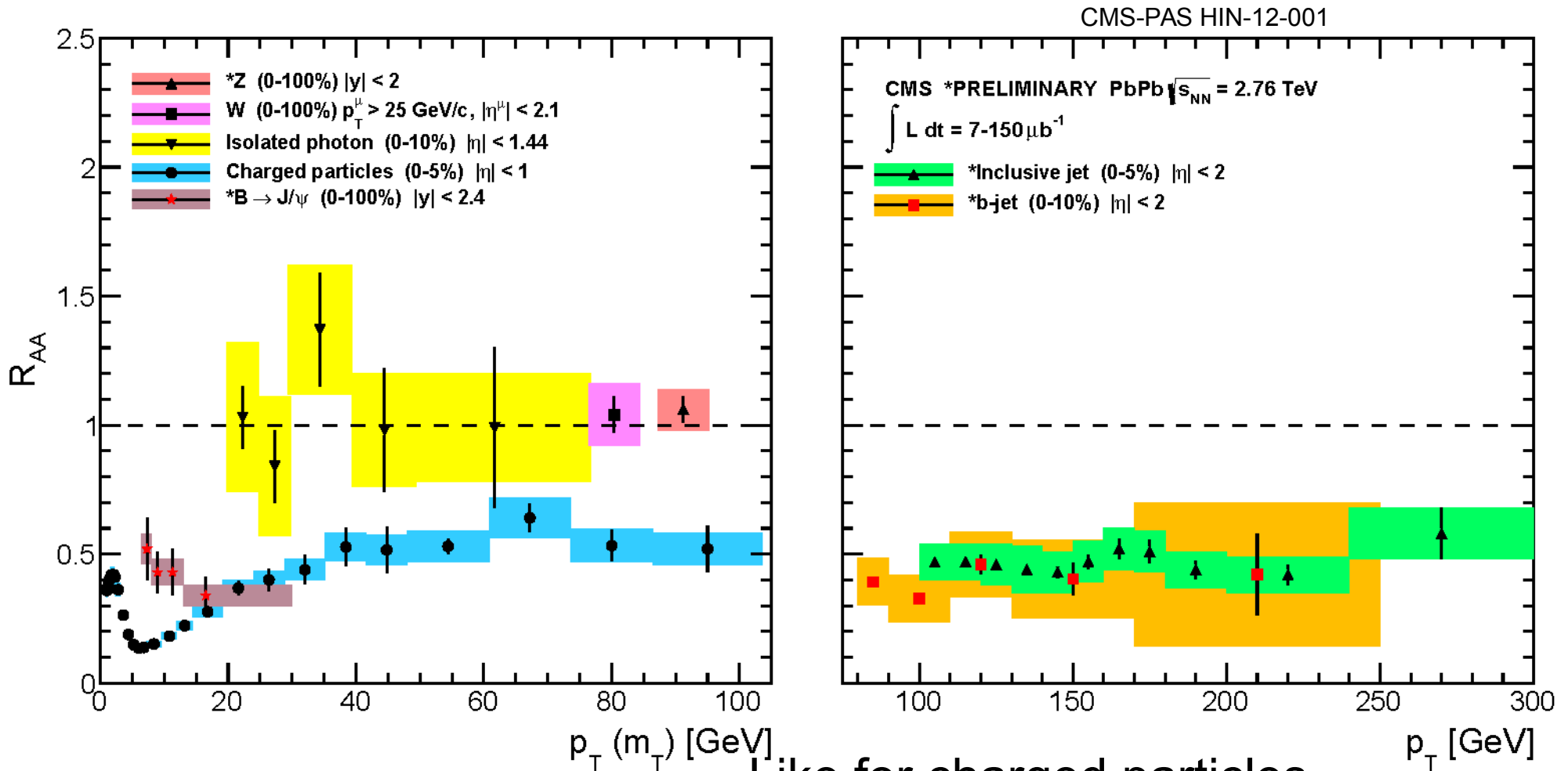
RHIC *LHC*



- *Electromagnetic probes* – consistent with no modification – medium is transparent to them
- *Strong probes* – significant suppression – medium is opaque to them

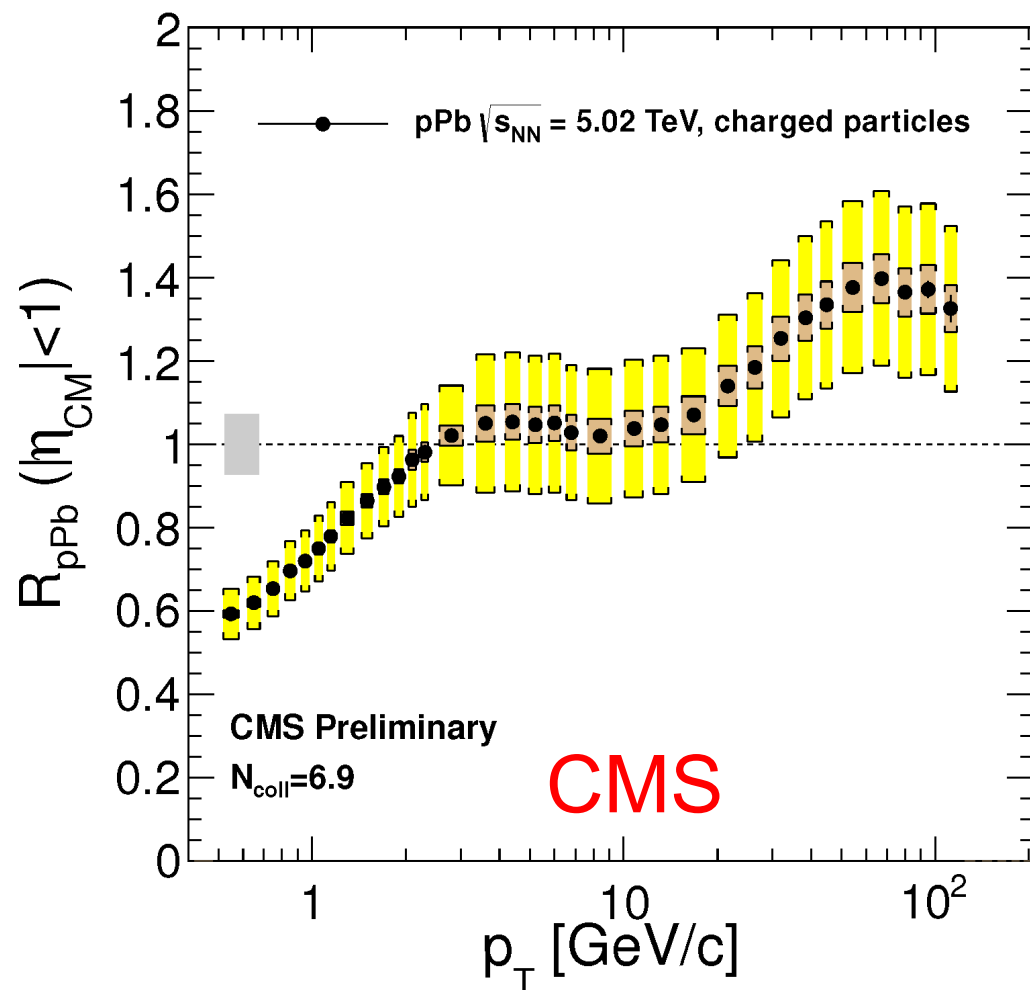
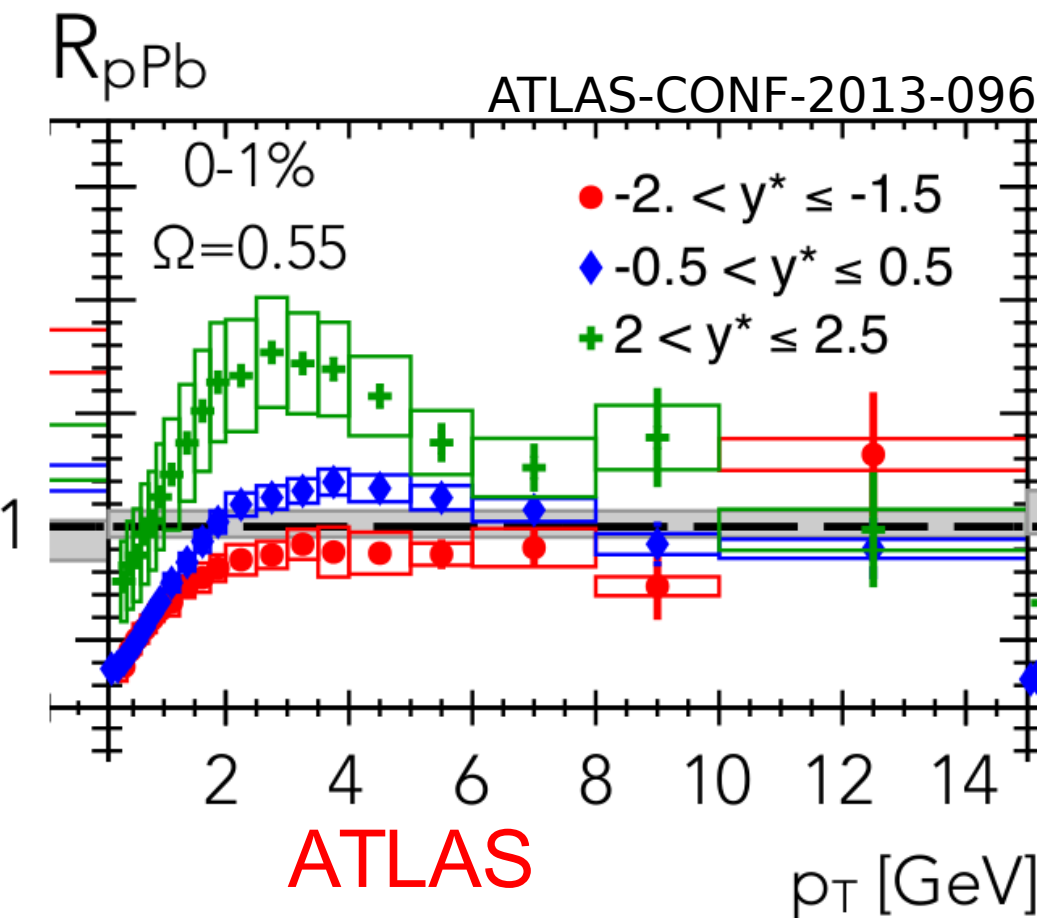
Nuclear modification factor R_{AA} at LHC

Fully unfolded inclusive jet R_{AA}
pp 2.76 TeV reference

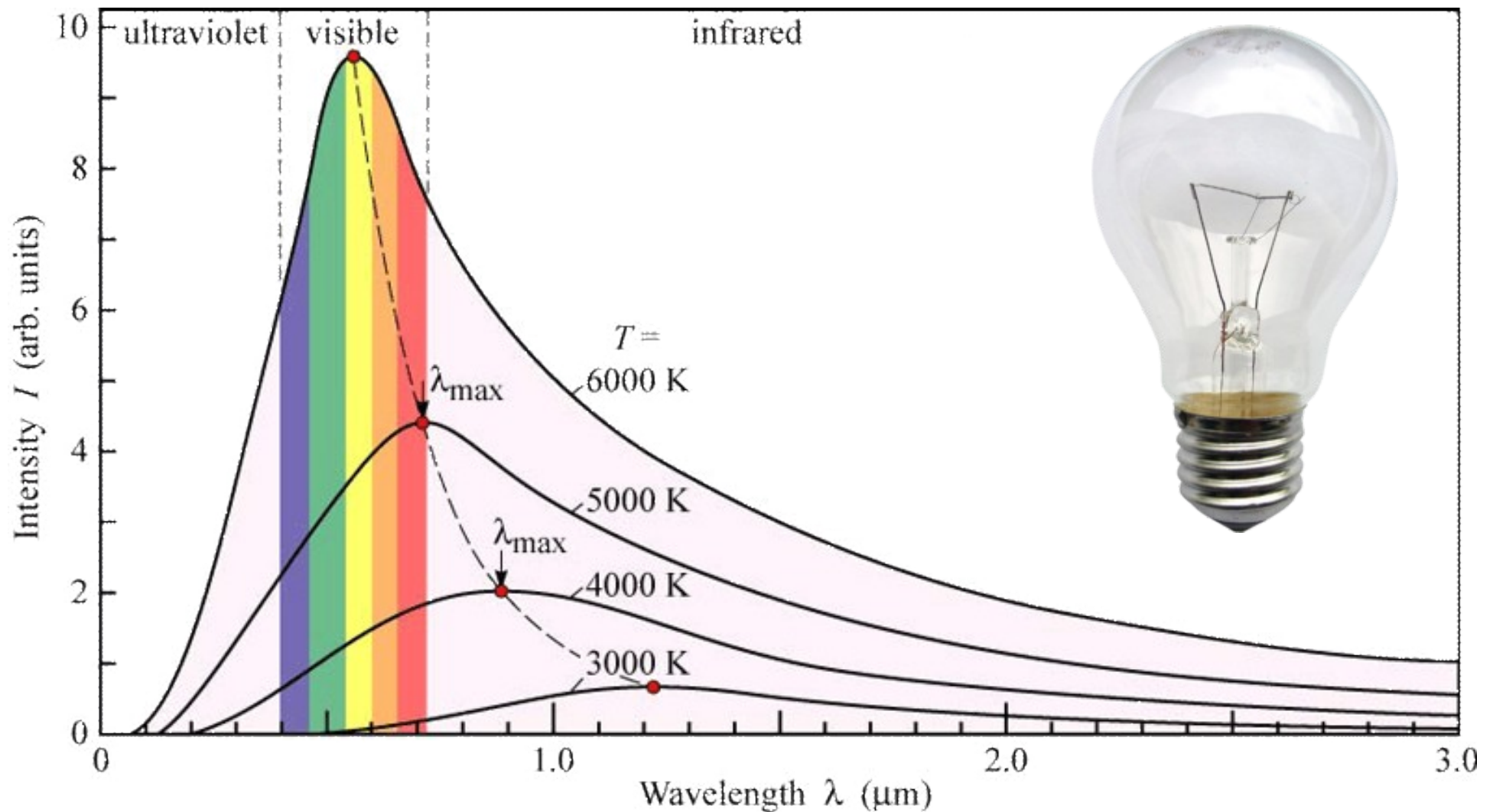


Like for charged particles,
high- p_T jet R_{AA} flat at ≈ 0.5

p+Pb as a control

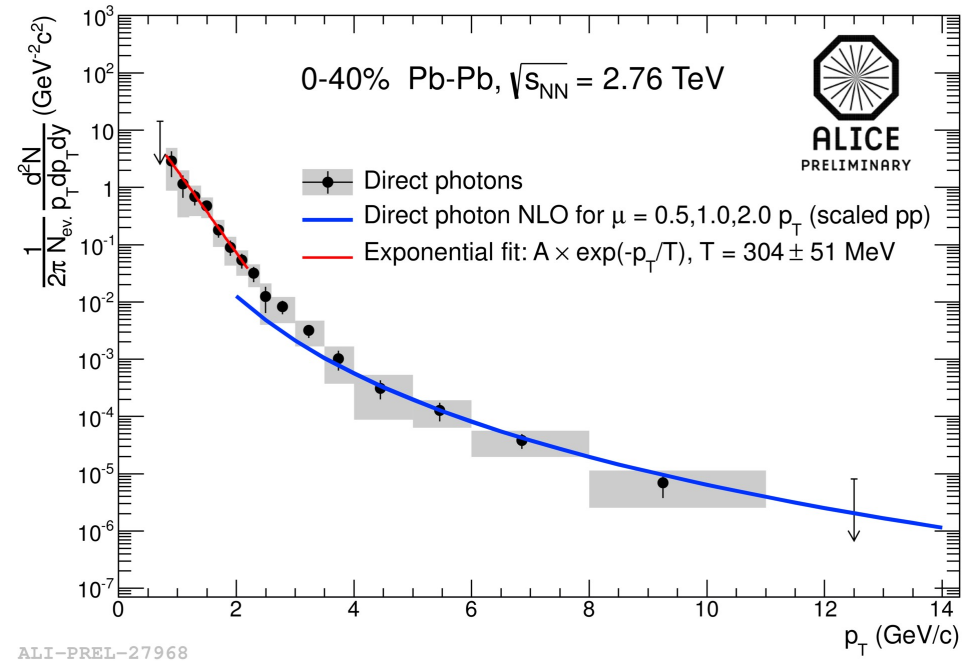
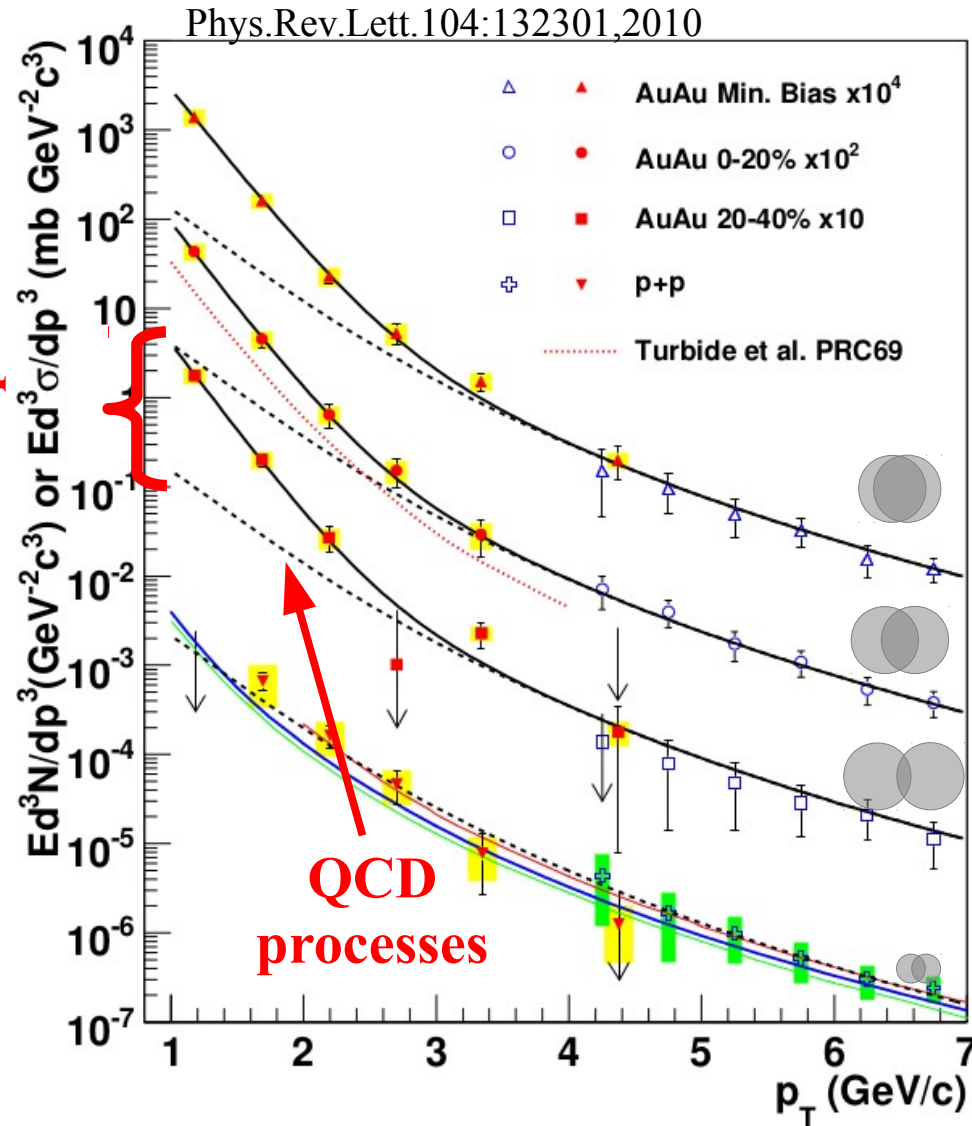


Measuring temperature



Thermal photons

Thermal photons

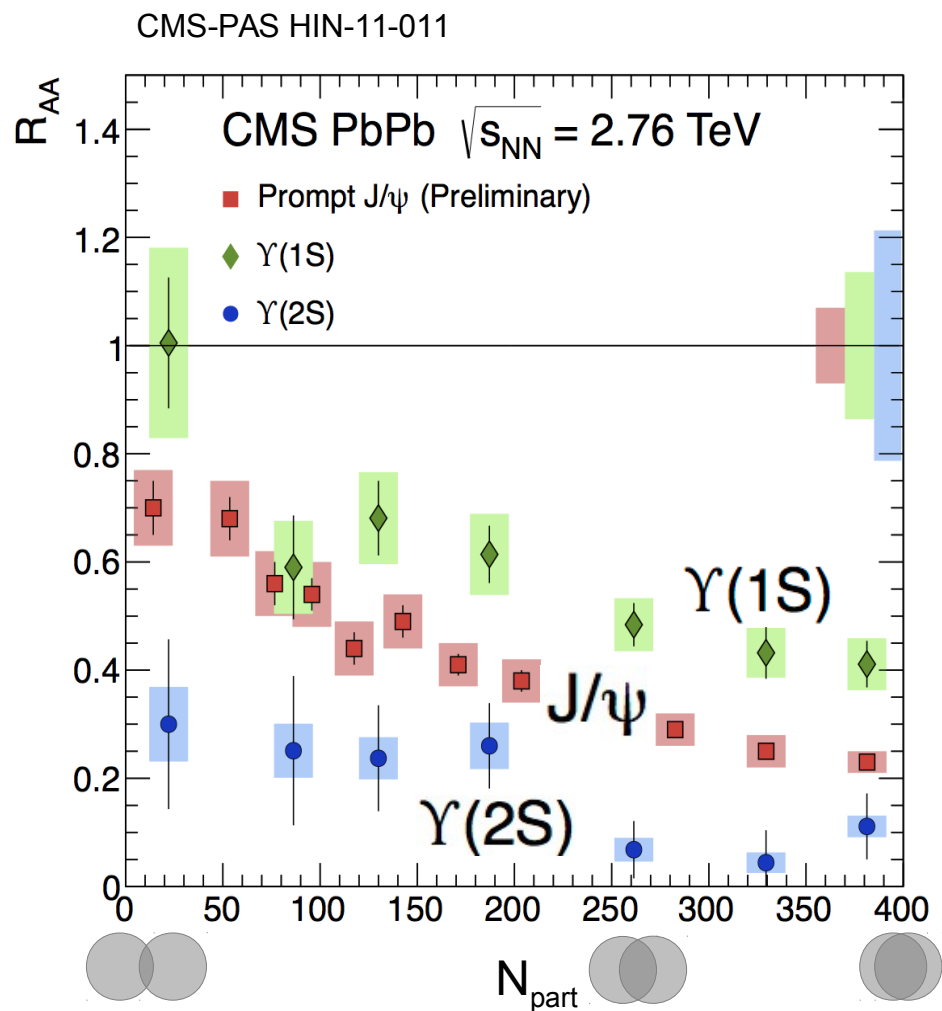


ALICE collaboration:
Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV
Inverse slope: $T = 304 \pm 51$

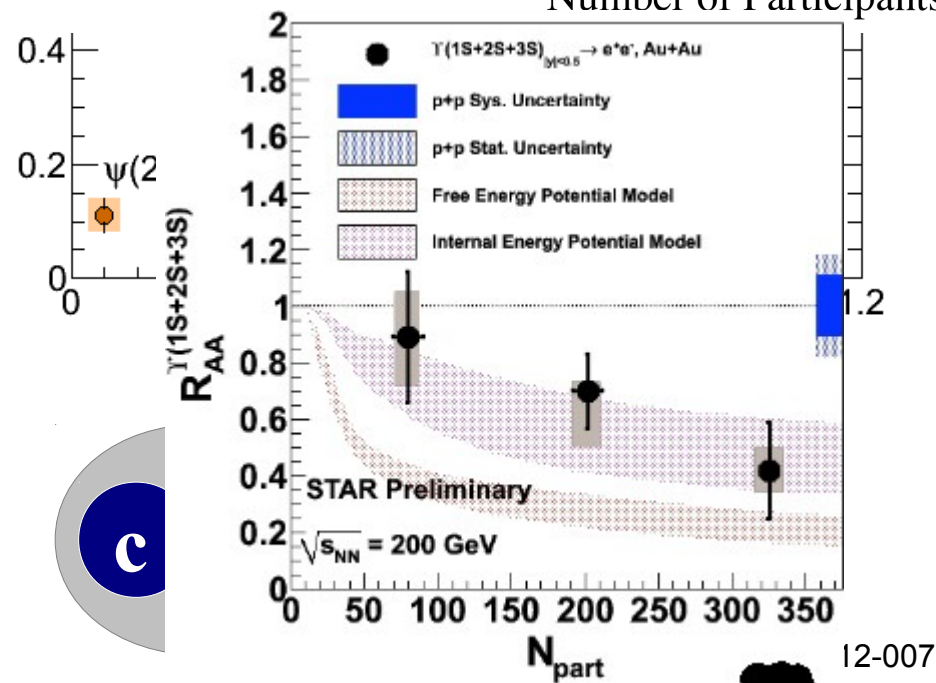
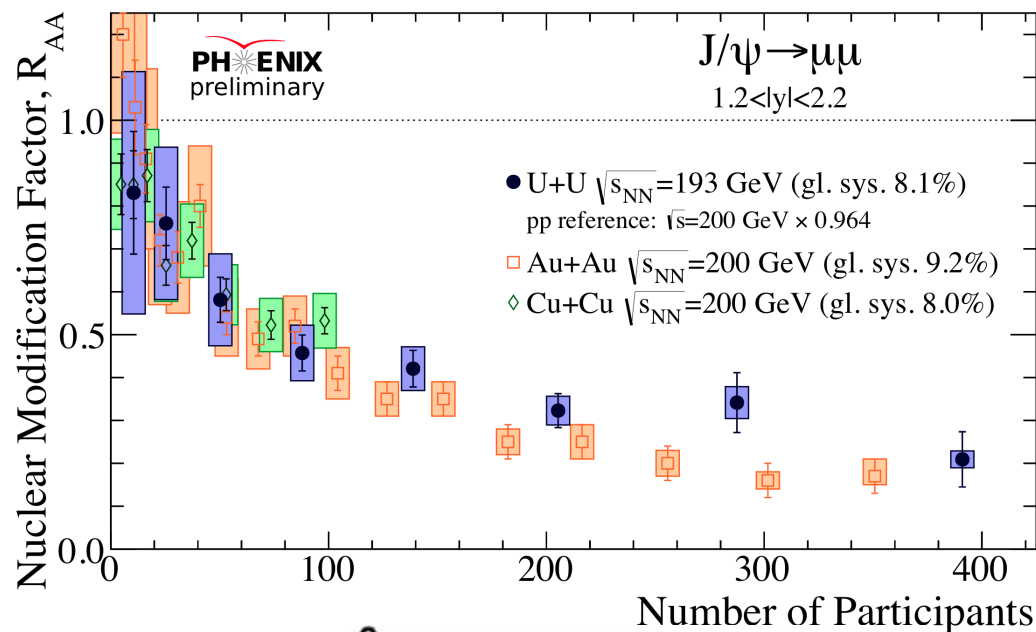
PHENIX collaboration: Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV

Inverse slope: $T = 221 \pm 19$ (stat) ± 19 (syst) MeV

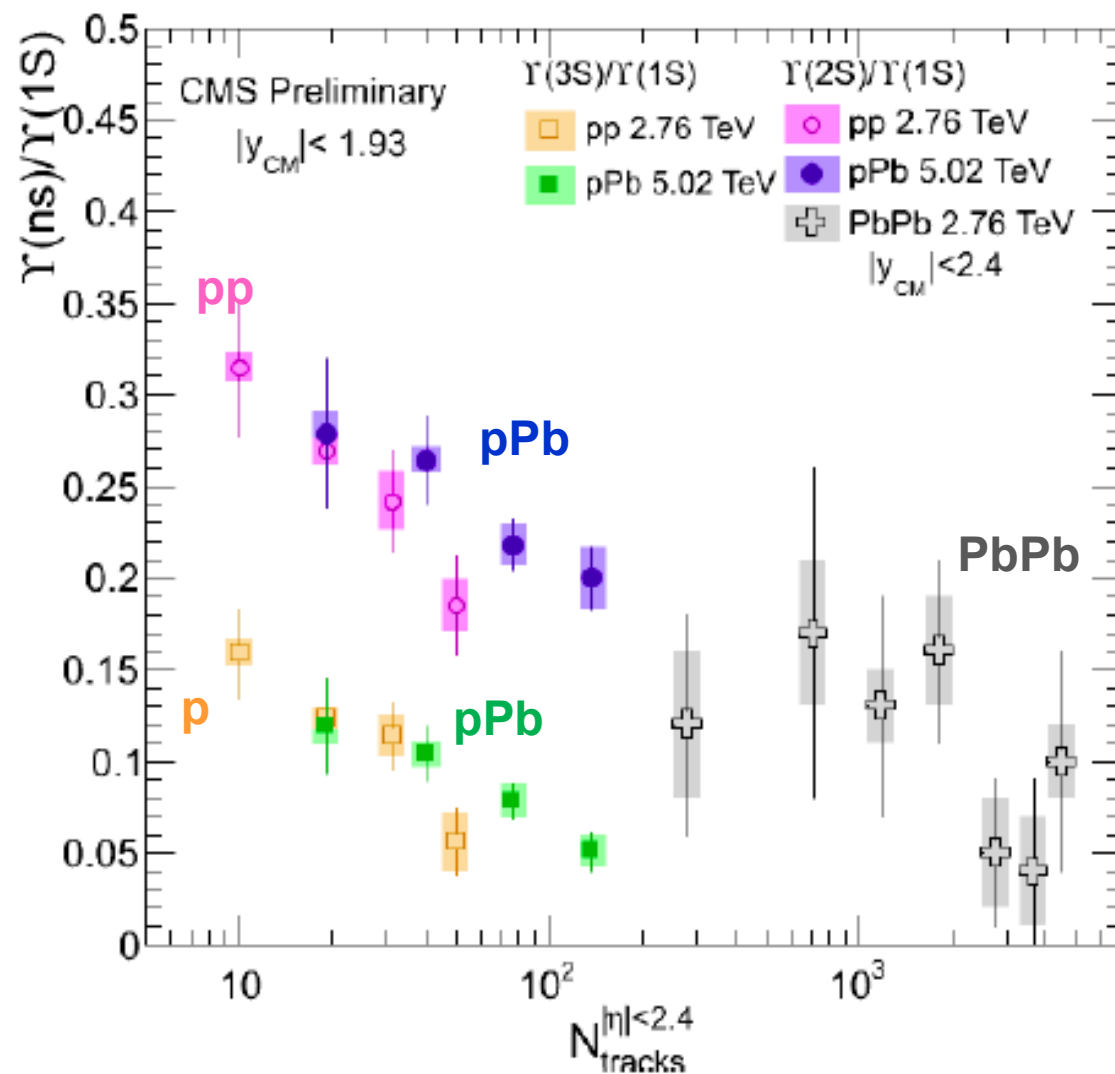
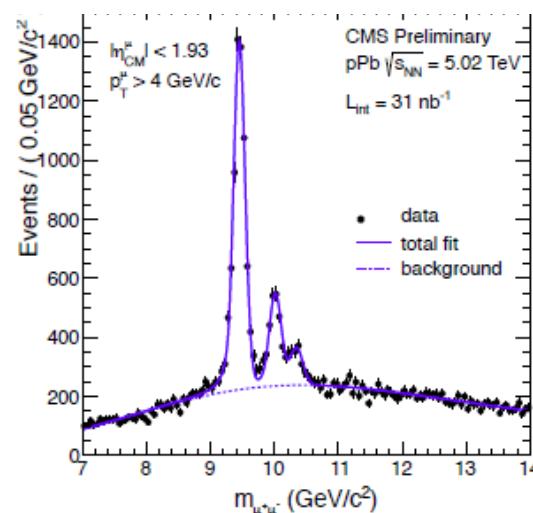
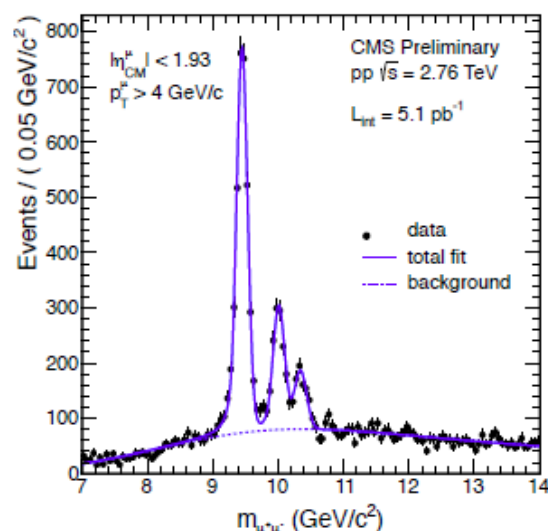
Building a quarkonium-thermometer



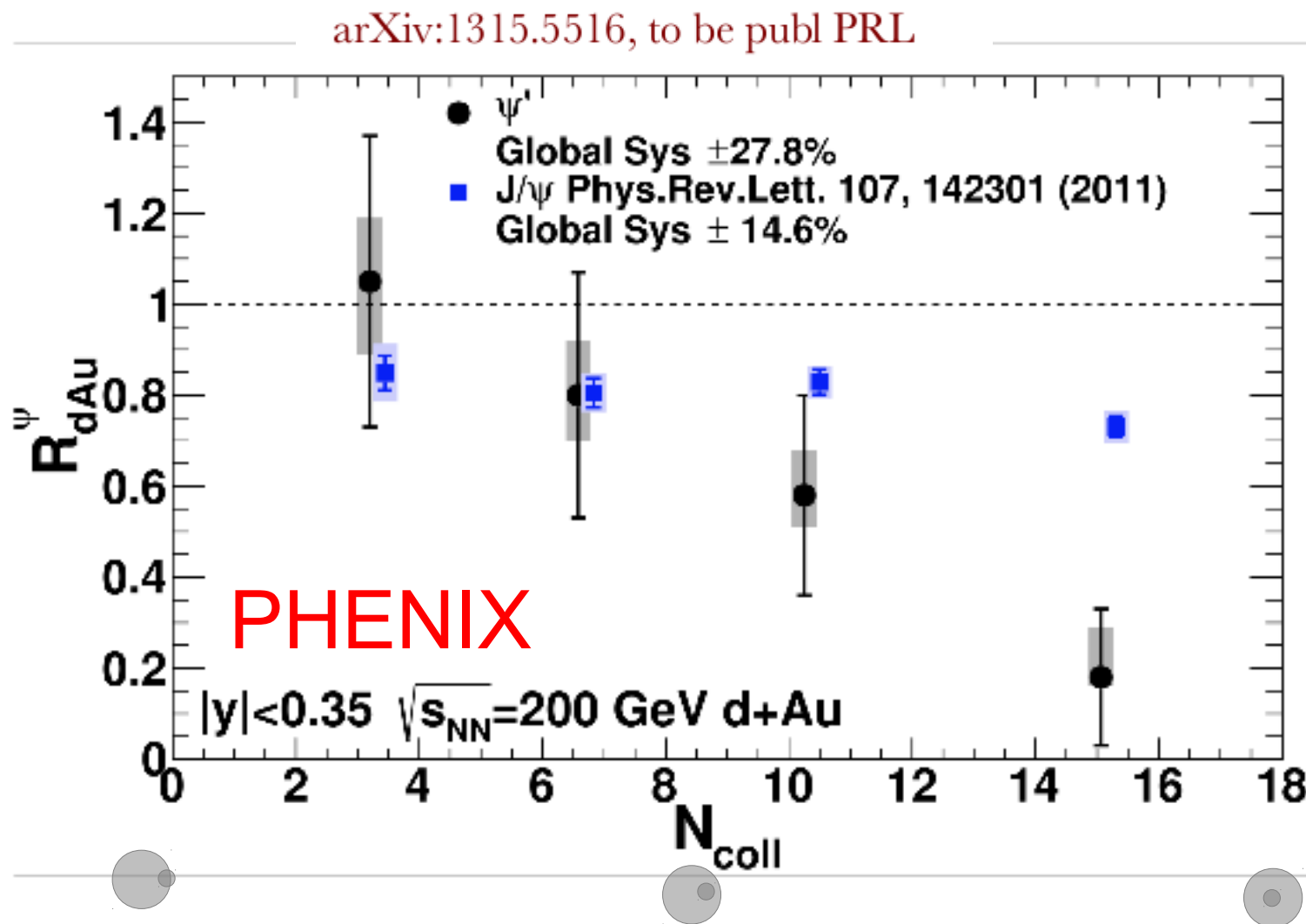
Clear hierarchy in R_{AA} of different quarkonium states



Suppression of quarkonia in p+Pb



Suppression of quarkonia in d+Au



Take home messages

- If we get nuclear matter dense enough, we make a new phase of matter, which we produce in high energy heavy ion collisions.
- This medium is transparent to colored probes and translucent to electromagnetic probes...
- ...And extremely hot and dense.

